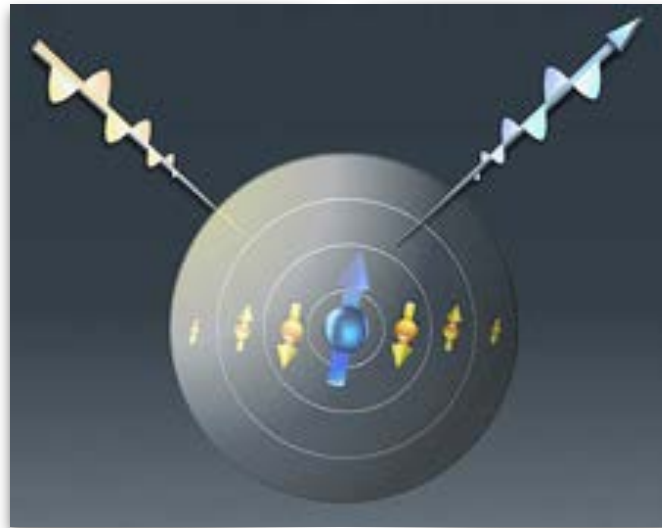
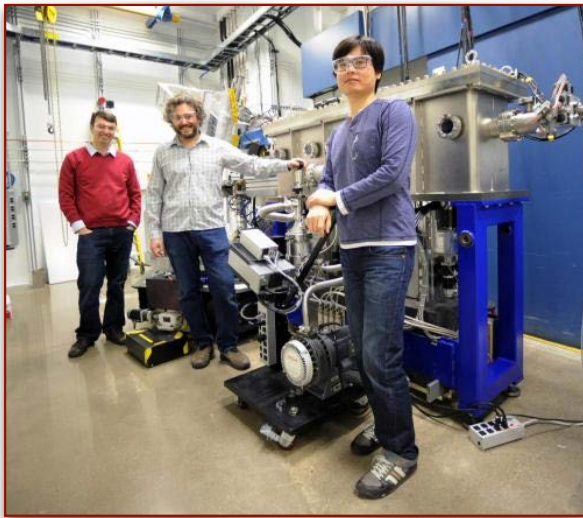


# Resonant Inelastic X-Ray Scattering & Electron Dynamics

Jason N. Hancock  
University of Connecticut  
jason.hancock@uconn.edu



Neutron and X-ray Scattering School 2019  
June 24, 2019, Argonne National Laboratory, Argonne, IL

# UConn Condensed Matter Physics



**Alexander Balatsky**  
Theory



**Gayanath Fernando**  
Theory



**Boris Sinkovic**  
Thin films synthesis  
Electron spectroscopy



**Elena Dormidontova**  
Soft Matter Theory



**Jason Hancock**  
THz, Infrared, X-ray  
Applied physics



**Ilya Sochnikov**  
Transport  
Scanning SQUID



**Niloy Dutta**  
Photonics  
Applied physics



**Menka Jain**  
Thin film synthesis



**Barrett Wells**  
PLD films, Muons,  
Neutrons, ARPES

---

Plus strong connections with: AMO group, UConn Institute for Materials Science, New UConn Tech Park

# Photon Science at UConn

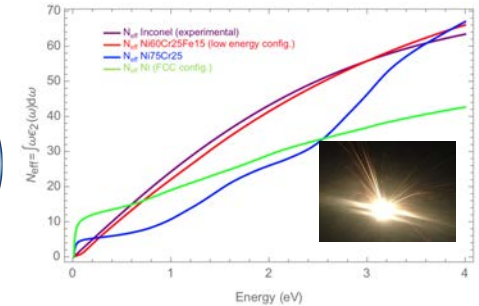
- Additive manufacturing



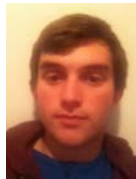
Erin Curry  
PhD student



Kaitlin Lyszak  
PhD student



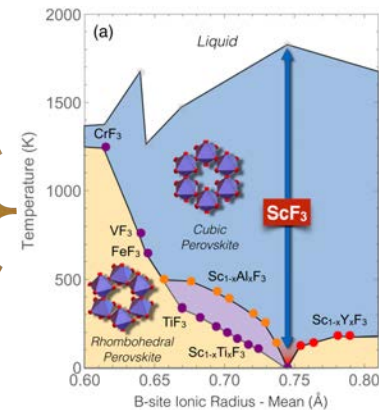
- Negative thermal expansion near structural quantum phase transitions



Connor Occhialini BS '18  
(PhD @ MIT Aug 2018)



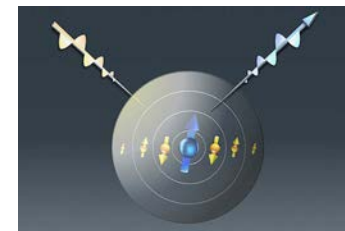
Sahan Handunkanda (PhD '18)



- *f*-electron physics via RIXS



Donal Sheets  
PhD student



My PhD research does/will likely use...



A

X-rays



B

Neutrons



C

Both X-rays and  
Neutrons



D

Neither



E

I am a theorist

My PhD research does/will likely regard...

**A**

**Mostly structural studies**

**B**

**Spectroscopy and electronic structure**

**C**

**Advanced imaging**

**D**

**Still figuring it out**

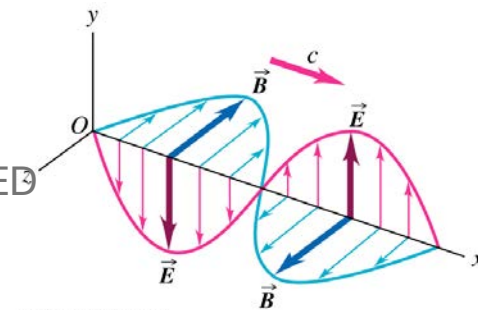
**E**

**None of these**

# **Spectroscopy in a nutshell**

# The Electromagnetic Field

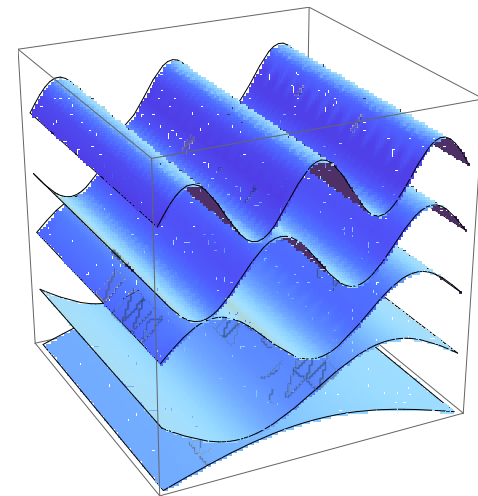
- Photons are bosons, independent, luminal
- Described by Maxwell's equations at low energy, QED at high energy
- Fully understood, quantized, modes are described as



$$|v\rangle = |n_0, \dots, n_{\mathbf{k}} \dots\rangle$$

Vector potential:

$$\mathbf{A}(\mathbf{r}, t) = \sum_{\mathbf{k}} \sum_{\mu=-1,1} \left( \mathbf{e}^{(\mu)}(\mathbf{k}) a_{\mathbf{k}}^{(\mu)}(t) e^{i\mathbf{k}\cdot\mathbf{r}} + \bar{\mathbf{e}}^{(\mu)}(\mathbf{k}) \bar{a}_{\mathbf{k}}^{(\mu)}(t) e^{-i\mathbf{k}\cdot\mathbf{r}} \right)$$

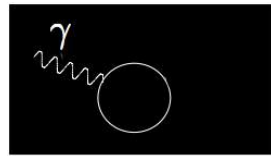


# Spectroscopy, general remarks

- Sample can exchange energy with EM field

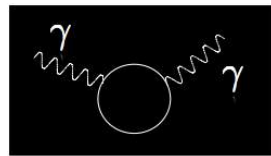
$$|\psi_g\rangle \otimes |v\rangle \rightarrow |\psi_e\rangle \otimes |v'\rangle$$

(sample)      (vacuum)



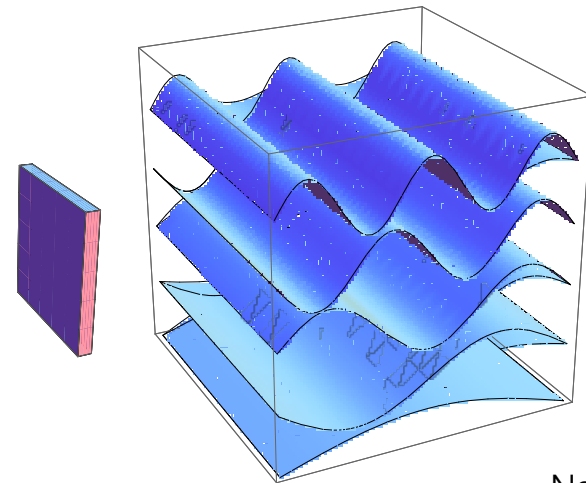
**One photon**  
Optical spectroscopy  
X-ray absorption

$$|n_0, \dots, n_{\mathbf{k}} \dots\rangle \rightarrow |n_0, \dots, n_{\mathbf{k}} - 1 \dots\rangle$$



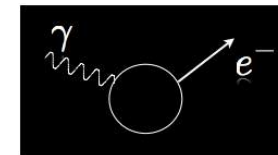
**Two photon**  
Raman  
Inelastic X-ray Scattering

$$|n_0, \dots, n_{\mathbf{k}} \dots\rangle \rightarrow |n_0, \dots, n_{\mathbf{k}} - 1 \dots, n_{\mathbf{k}'} + 1 \dots\rangle$$

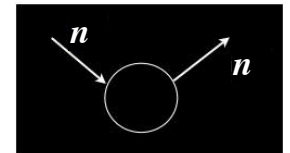


Related spectroscopies:

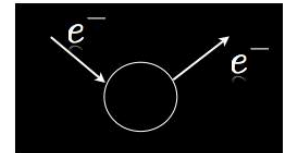
ARPES



Neutron



EELS

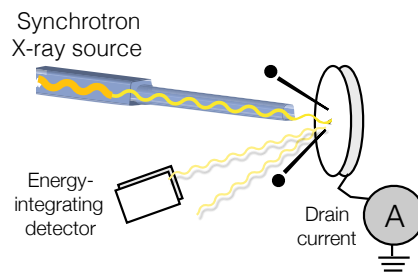
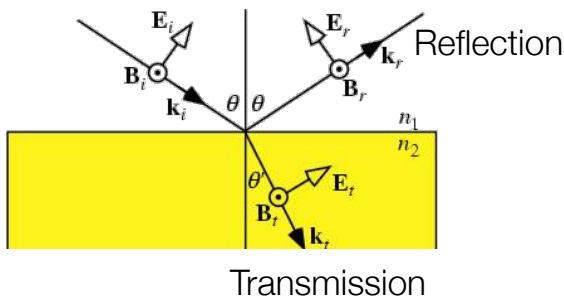




# Photon spectroscopies

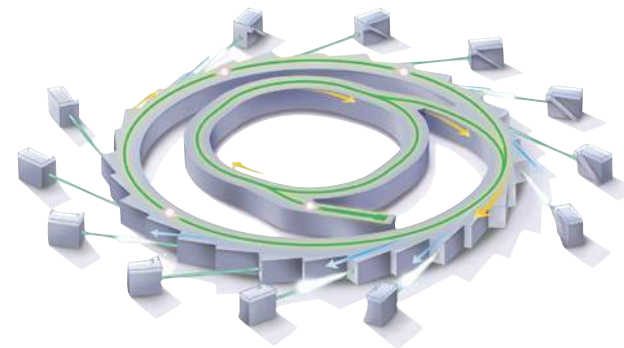
## One-photon

- Can measure from reflection or transmission in time or frequency domain
- Use Fresnel's equations to determine  $n + ik$ ,  $\sigma_1 - i\sigma_2$ ,  $\epsilon_1 + i\epsilon_2$  and relate to fundamental behavior
- X-ray absorption collected differently, through total electron or total fluorescence yield



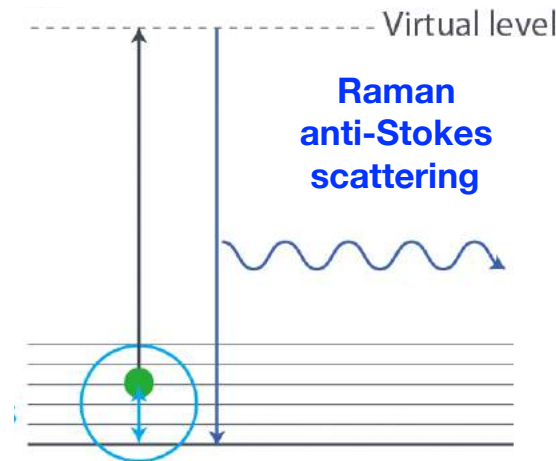
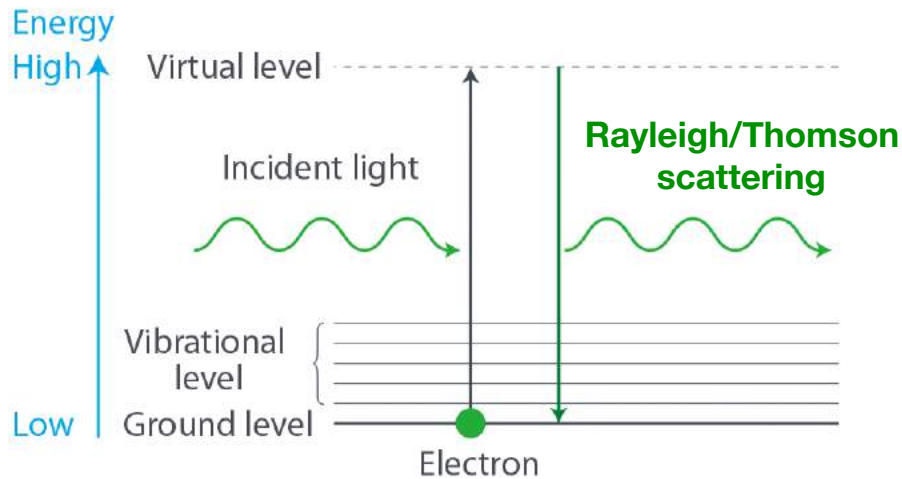
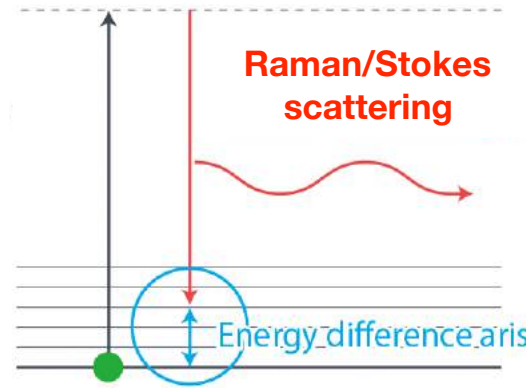
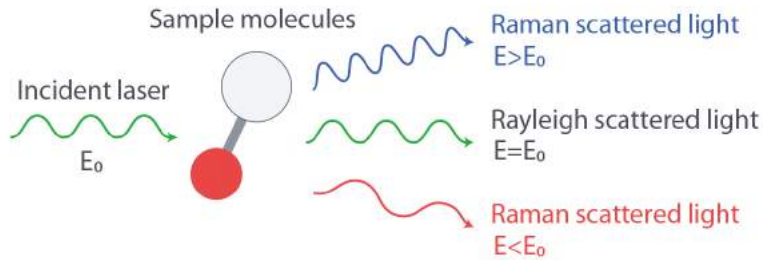
## Two-photon

- IXS and RIXS (also REXS, Raman)
- Probes electronic excitations in a momentum-resolved way
- Technical advances making rapid progress



# Two-photon scattering examples

# Raman=inelastic light scattering



# Modern Raman scattering



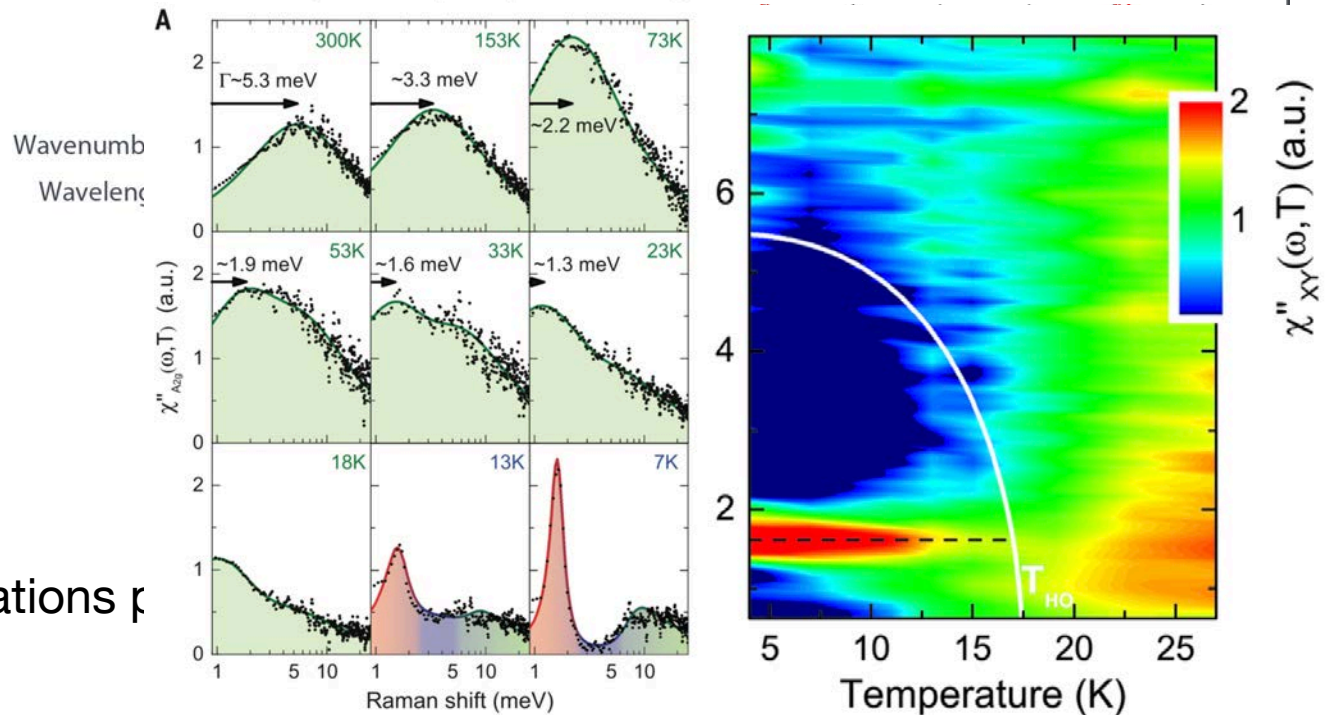
- Laser source=high resolution
- Lattice excitations easy to see
- Electronic and magnetic excitations  $\chi''$

HEAVY FERMIONS

Kung, et al Science 347, 1339 (2015)

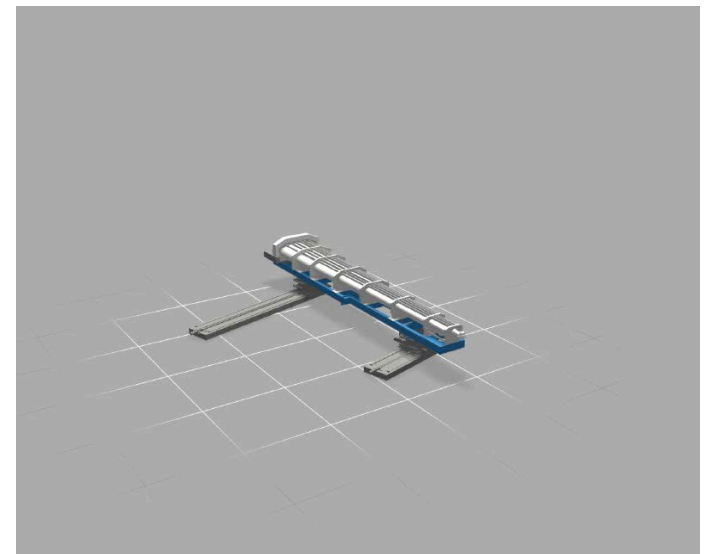
## Chirality density wave of the “hidden order” phase in URu<sub>2</sub>Si<sub>2</sub>

H.-H. Kung,<sup>1\*</sup> R. E. Baumbach,<sup>2†</sup> E. D. Bauer,<sup>2</sup> V. K. Thorsmølle,<sup>1‡</sup> W.-L. Zhang,<sup>1</sup>  
K. Haule,<sup>1\*</sup> J. A. Mydosh,<sup>3</sup> G. Blumberg<sup>1,4\*</sup>

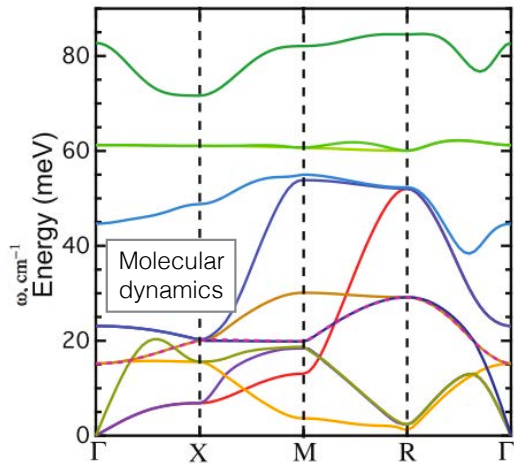
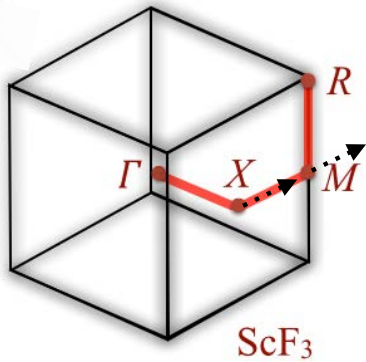


# High Energy Resolution Inelastic X-ray Scattering

- Sector 30, Advanced Photon Source, Argonne National Lab
- 23,724 eV incident energy
- <1 meV incident bandwidth
- Resolving power  $E_i/\Delta E_i = 2 \times 10^7$
- 9 analyzers sample, 9 momenta transfers simultaneously, 9m arm
- 1.5 meV energy resolution
- 20 $\mu\text{m}$  x 5 $\mu\text{m}$  spot size
- Measures energy and momentum distribution of lattice vibrations  $S(\vec{q}, \omega)$

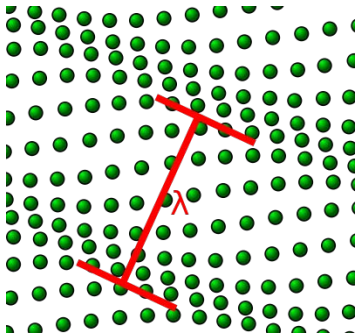


# IXS from phonons



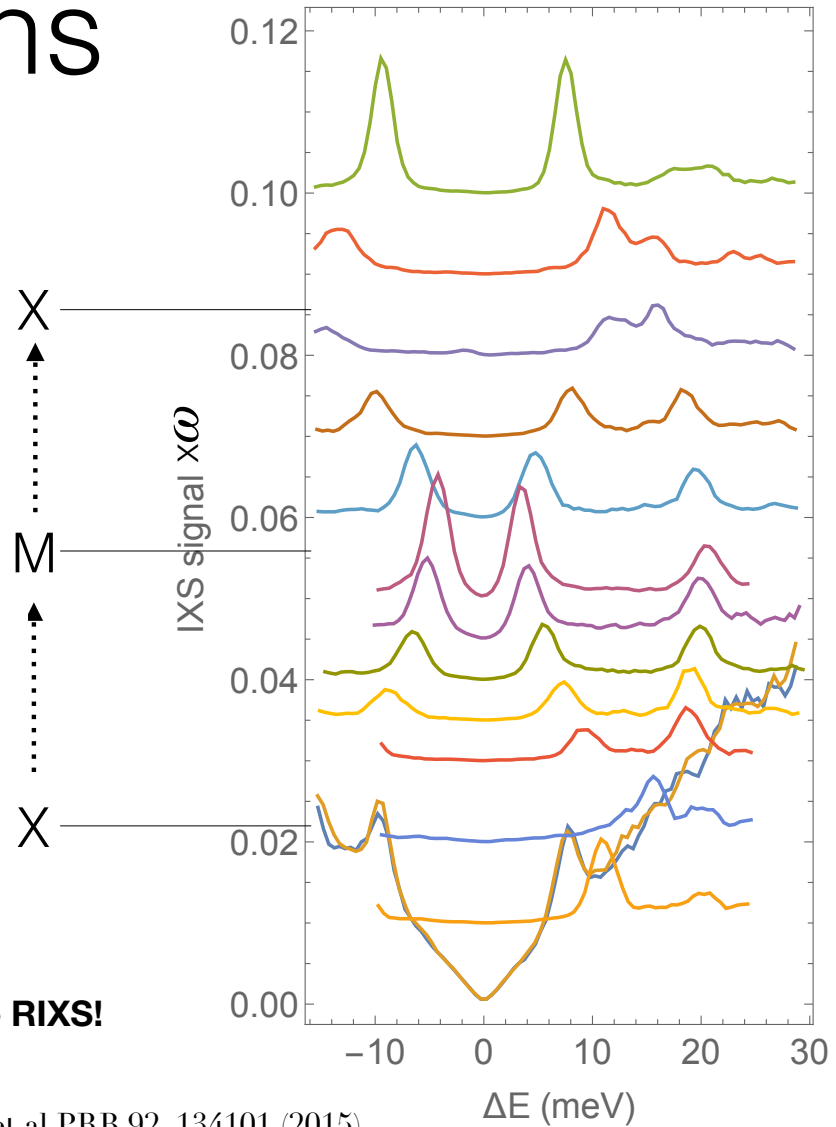
Li et al, PRL 107 195504 (2011)

- IXS measures collective modes of lattice system: phonons



Energy of each mode at wavevector  $k$

What about electronic or magnetic collective modes? Use RIXS!



S. U. Handunkanda, et al PRB 92, 134101 (2015)

# What is this?



The emission spectrum of a star at  $z=0$  redshift



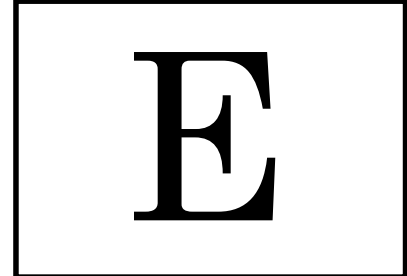
A flat rainbow snake



A fashionable belt

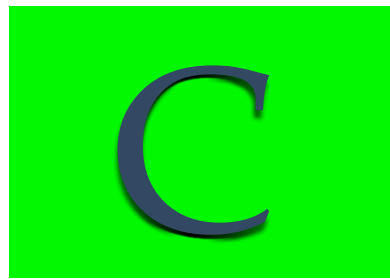
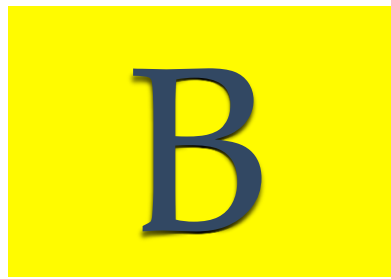


Balmer series of emission by hydrogen



There is more than one correct answer

# What is this?



The absorption spectrum of a star at  $z=0$  redshift

A flat rainbow snake

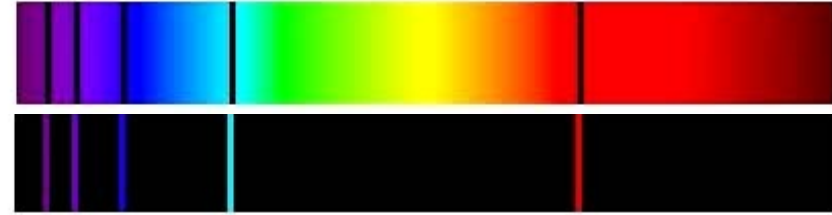
A fashionable belt

Balmer series of absorption by hydrogen

There is more than one correct answer

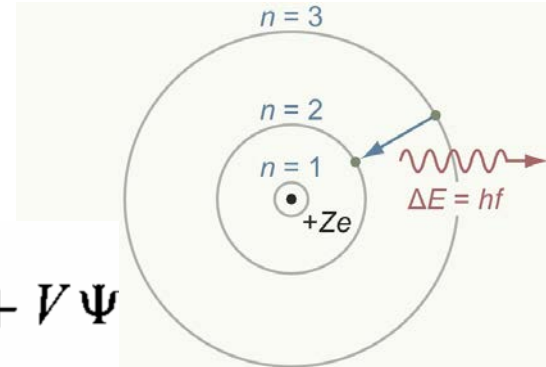


# Balmer series



- Describes transitions between quantized energy levels in hydrogen (one proton, one electron)

$$\frac{1}{\lambda} = R_H \left( \frac{1}{2^2} - \frac{1}{n^2} \right)$$

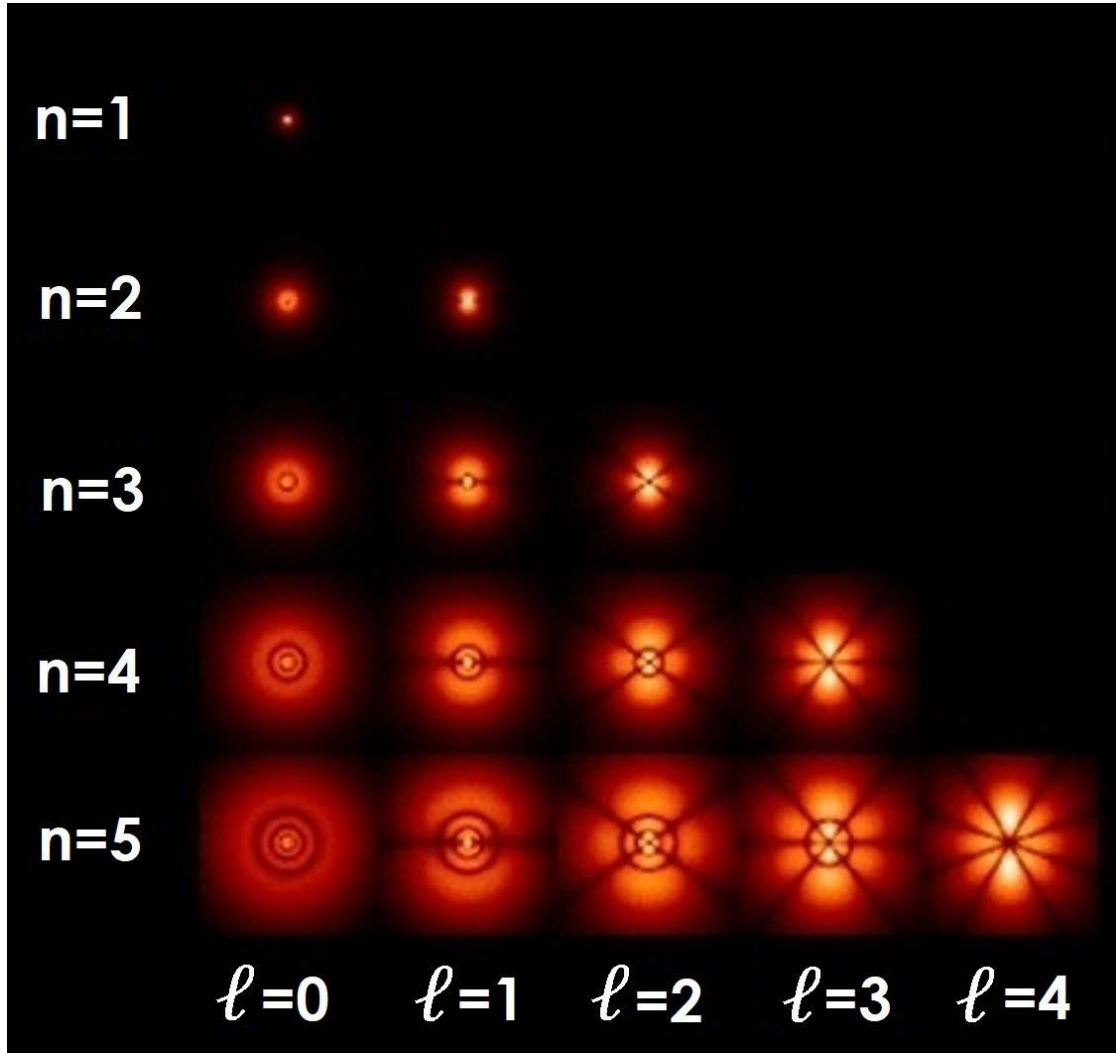


Schrodinger equation for 1/r potential: 
$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \Psi + V \Psi$$

Hydrogen wavefunctions: 
$$\psi_{nlm} = \sqrt{\left(\frac{2}{na}\right)^3 \frac{(n-l-1)!}{2n[(n+l)!]^3}} e^{-r/na} \left(\frac{2r}{na}\right)^l L_{n-l-1}^{2l+1} \left(\frac{2r}{na}\right) Y_l^m(\theta, \phi)$$

Energy levels: 
$$E_n = - \left[ \frac{m}{2\hbar^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \right] \frac{1}{n^2}$$

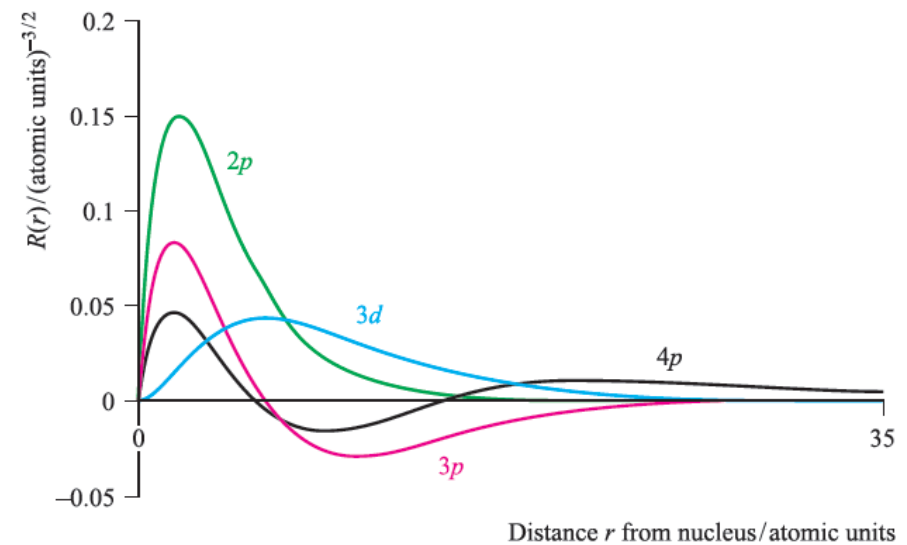
Bohr radius: 
$$a \equiv \frac{4\pi\epsilon_0\hbar^2}{me^2} = 0.529 \times 10^{-10} \text{ m}$$



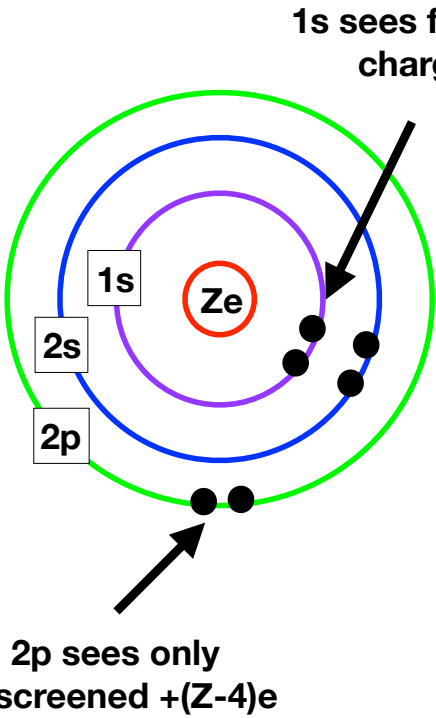
Quantum numbers:  $n, l, m$

Generalize for nuclear charge  $+Ze$ :

$$E_{z,n} = -\frac{Z^2 R_H}{n^2} \quad a_{Z,n} = \frac{na_0}{Z}$$



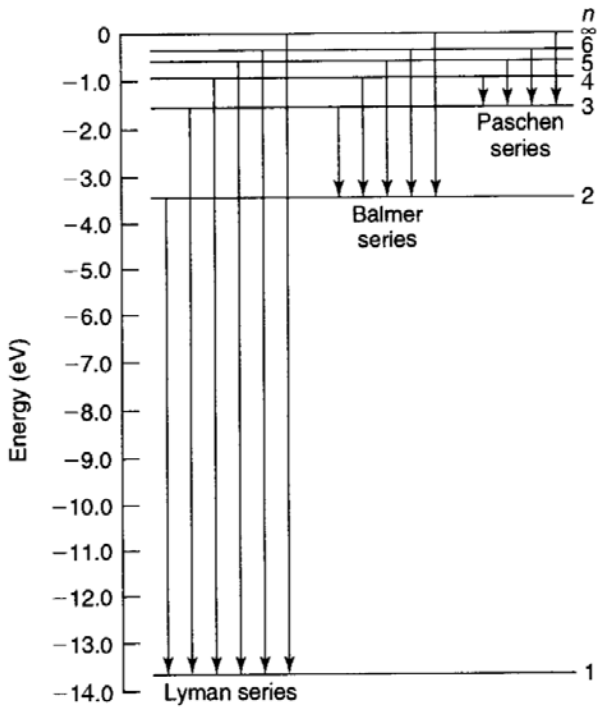
# How does Balmer series change multi-electron atoms like C?



Inner “core” electrons are much smaller than outer and can effectively screen nucleus for other electrons

Core electrons much more tightly bound than outer ones

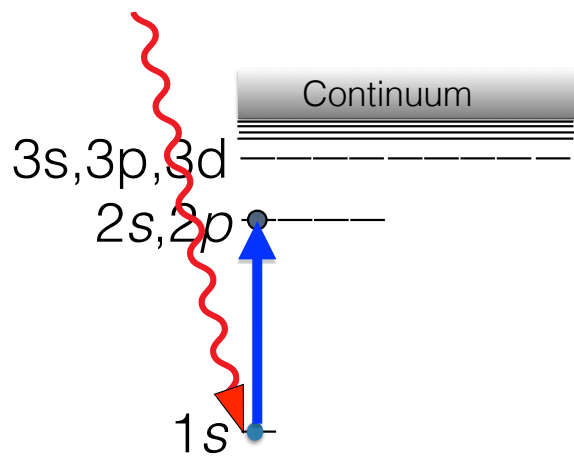
Lyman (1->n) series for multi-electron atoms form “K” X-ray edges



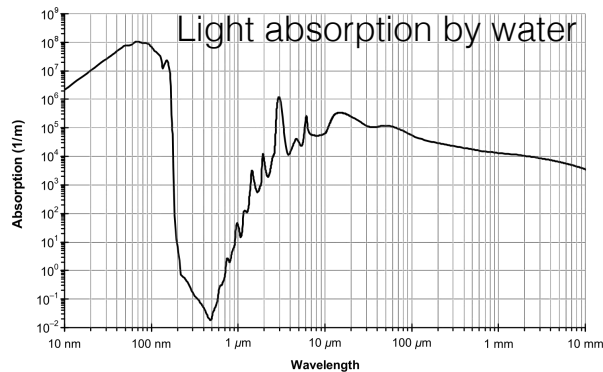
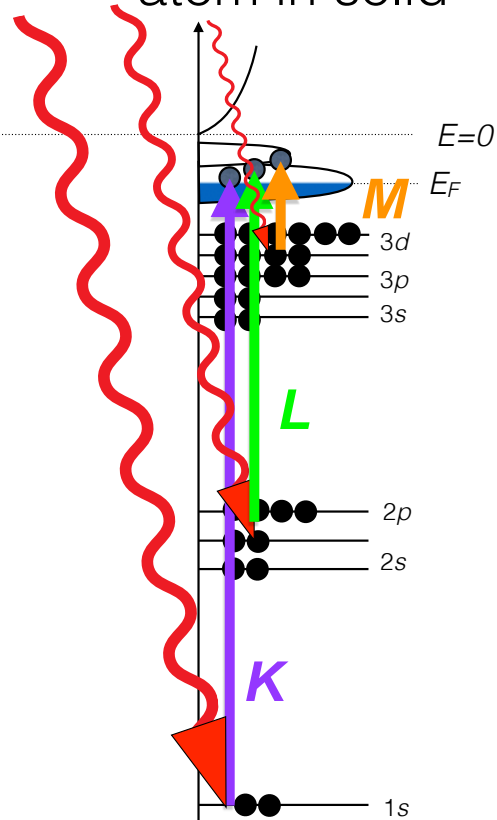
# Atomic transitions in solids



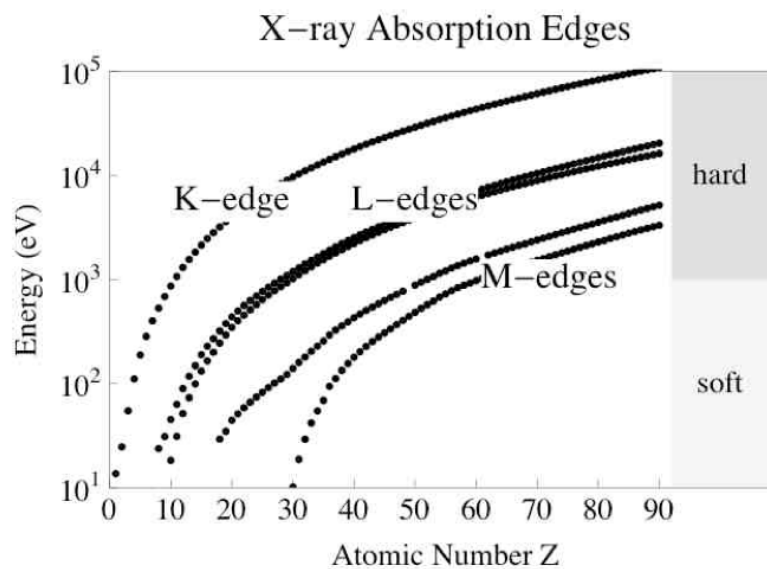
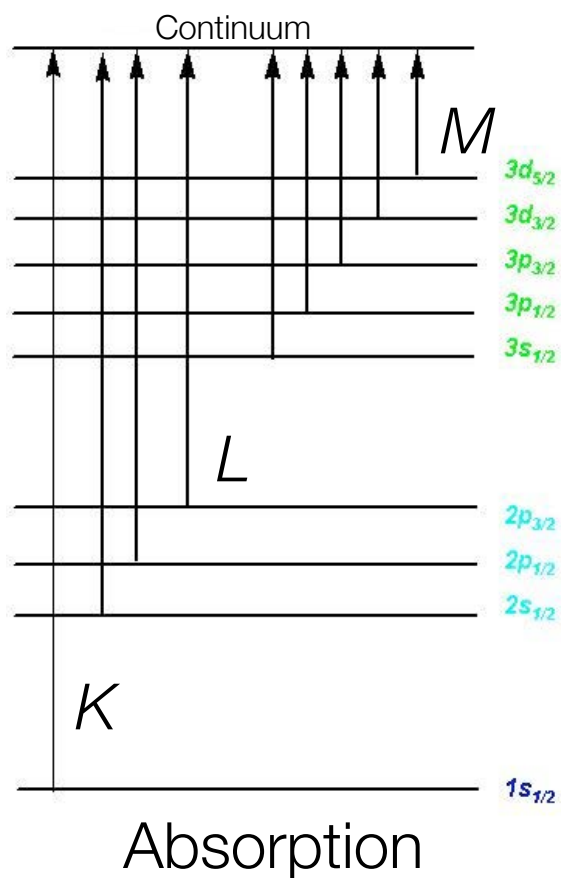
Hydrogen atom



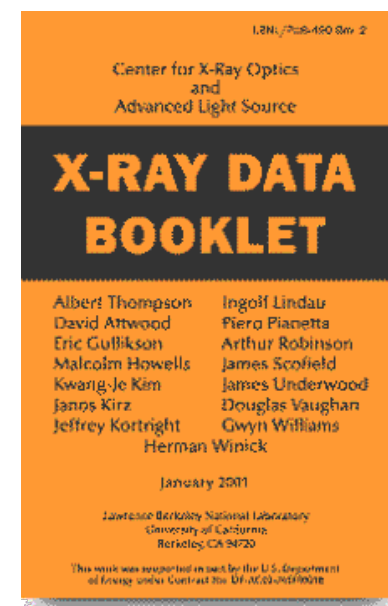
Multi-electron atom in solid



# X-ray transition lines

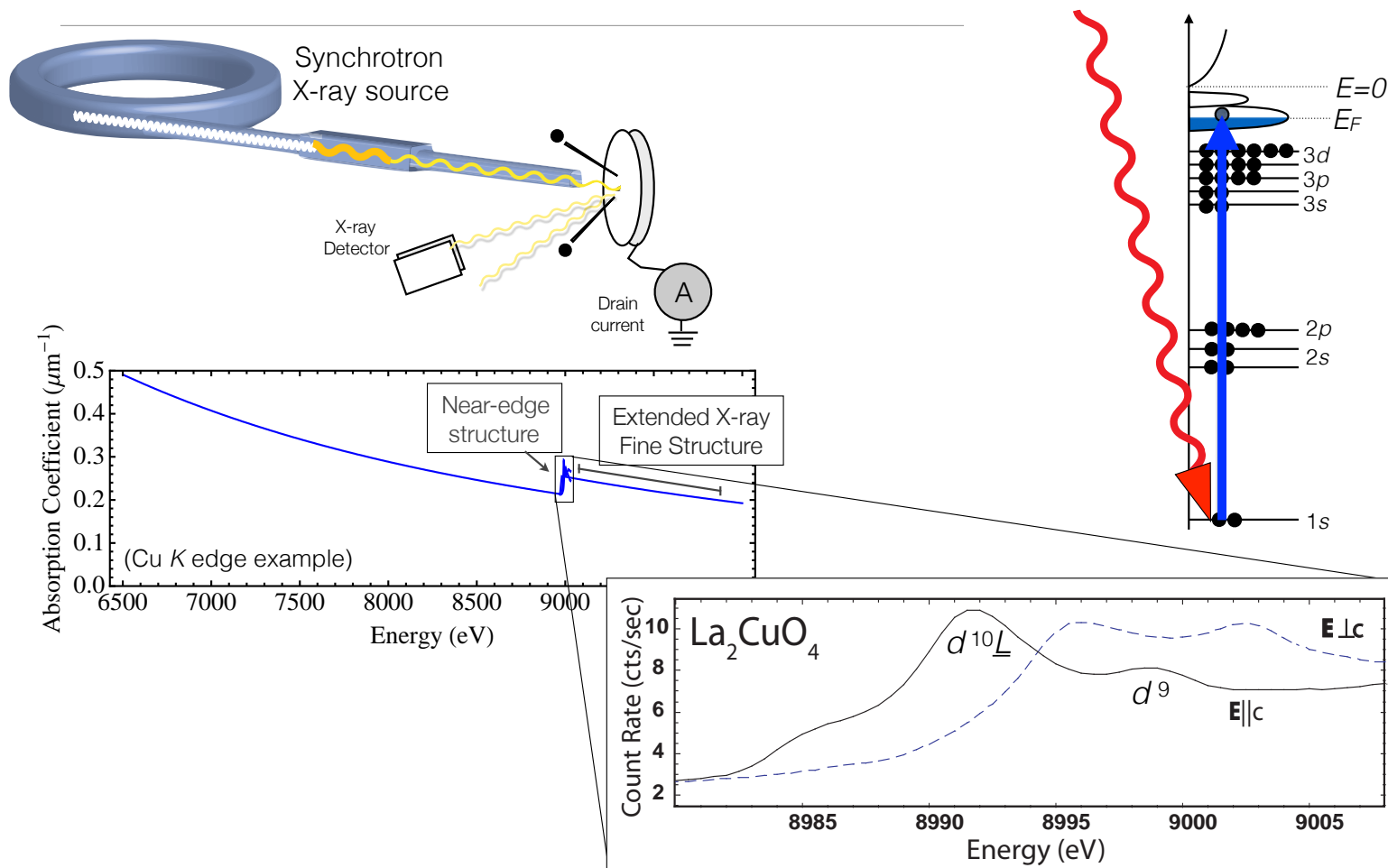


- Absorptive transitions of higher Z atoms lie in X-ray region

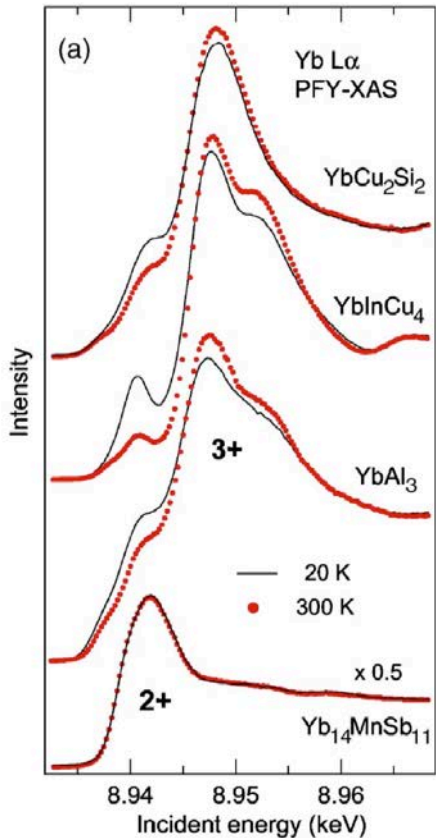


<http://xdb.lbl.gov>

# X-ray absorption spectroscopy (XAS)



# X-ray absorption spectroscopy, uses to determine valence state



## Example: Yb

Has two stable valence states

Yb<sup>+2</sup>, has  $f^{14}$  full shell - nonmagnetic ion

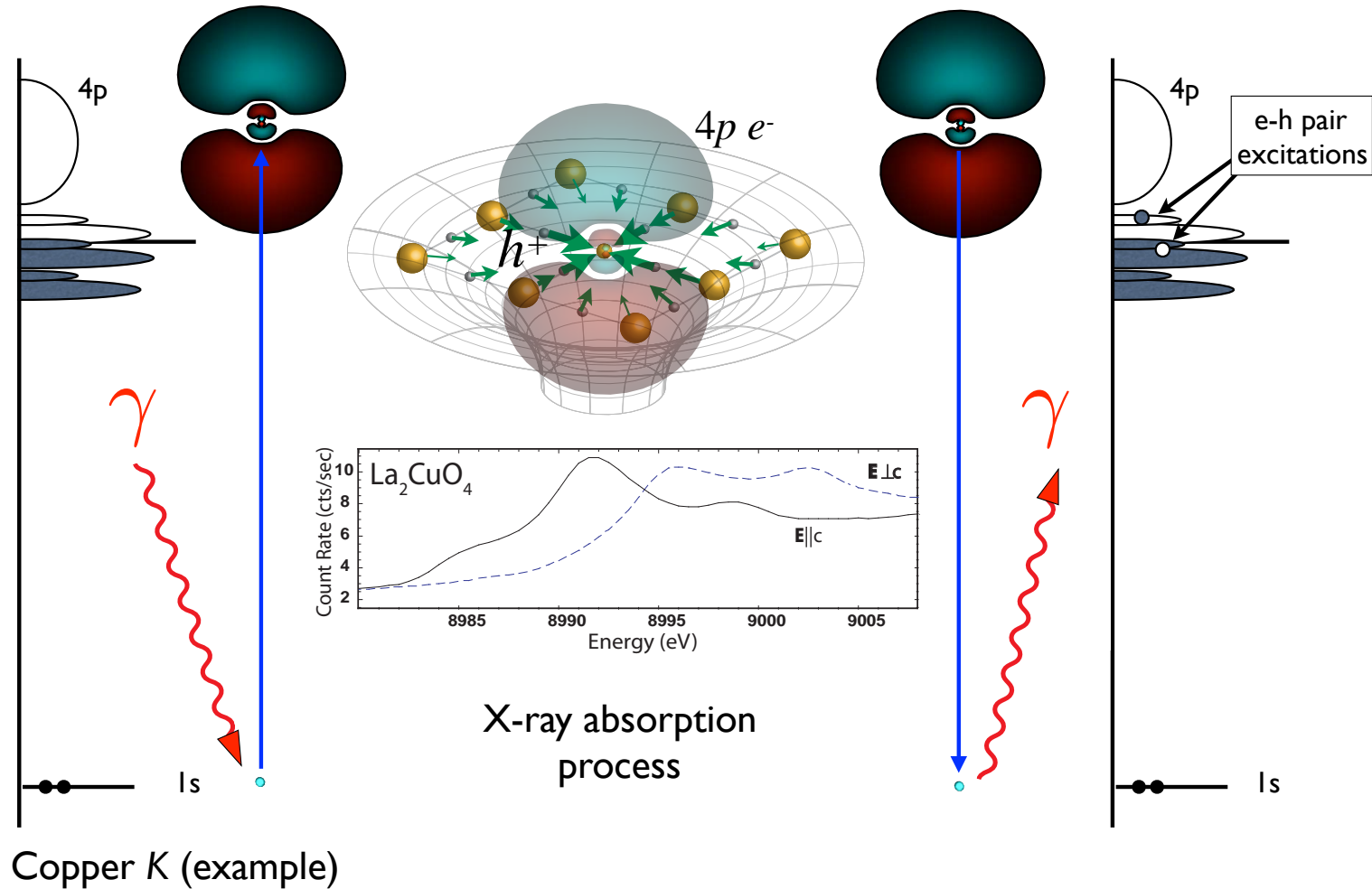
Yb<sup>+3</sup> has  $f^{13}$  one hole in f shell - long  $j=7/2$  mag moment

Real materials can have things in between, depends on interactions

Yb<sub>14</sub>MnSb<sub>11</sub> - no magnetic moment coming from Yb

YbAl<sub>3</sub> - long moment paramagnet from Yb<sup>+3</sup> in  $f^{13}$  state

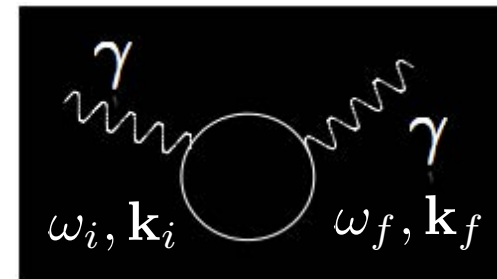
# X-ray edge absorption... and RIXS





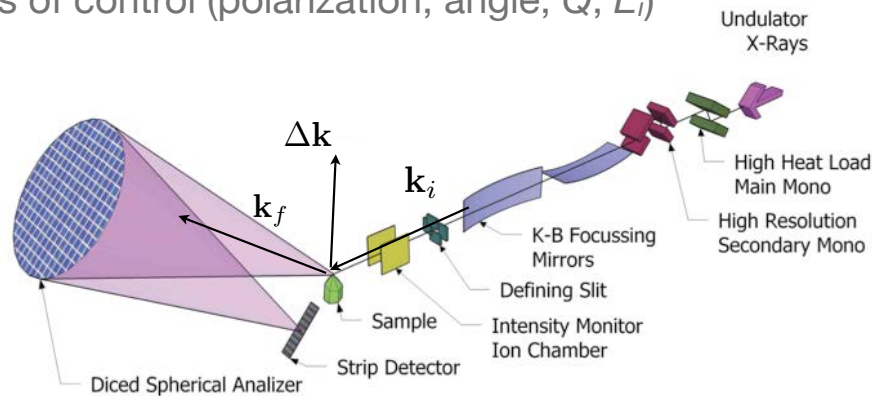
# Resonant Inelastic X-ray Scattering (RIXS)

- Resonant Raman spectroscopy, with resonance at an X-ray edge
- Described by Kramers-Heisenberg formula
- Large momentum transfer
- Atomic-species and valence-specific information
- Lots of control (polarization, angle,  $Q$ ,  $E_i$ )



$$\Delta\omega = \omega_i - \omega_f$$

$$\Delta\mathbf{k} = \mathbf{k}_i - \mathbf{k}_f$$



**Kramers-Heisenberg:**

$$\frac{d^2\sigma}{d\Omega_{k'} d(\hbar\omega'_k)} = \frac{\omega'_k}{\omega_k} \sum_{|f\rangle} \left| \sum_{|n\rangle} \frac{\langle f|T^\dagger|n\rangle \langle n|T|i\rangle}{E_i - E_n + \hbar\omega_k + i\frac{\Gamma_n}{2}} \right|^2 \delta(E_i - E_f + \hbar\omega_k - \hbar\omega'_k)$$

# RIXS today



**Resonance Raman Scattering**

Loss features in the RIXS intensity are represented on a plot of the scattered photon energy  $\hbar\omega'$  and incident photon energy  $\hbar\omega$  and incident momentum  $Q$  at zero transferred momentum. Emission lines appear as diagonal features in this  $\hbar\omega$ - $\hbar\omega'$  plot because in this case energy of the emitted photon  $\hbar\omega_k'$  is roughly constant. From Hancock *et al.*, 2009.

Raman losses

Emission line

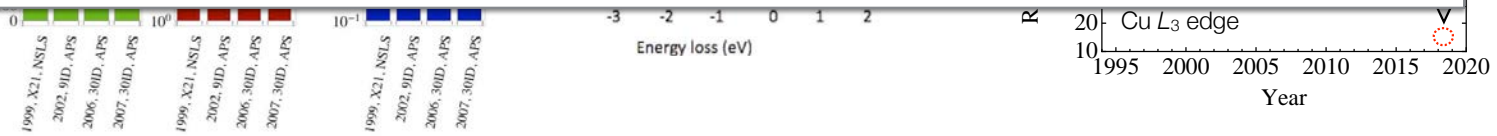
**Daniel Mazzone**

**Donal Sheets**

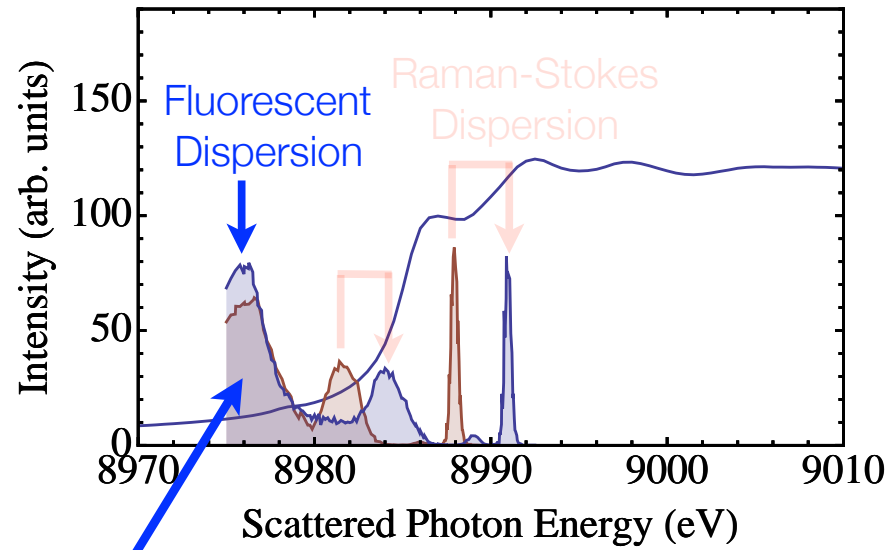
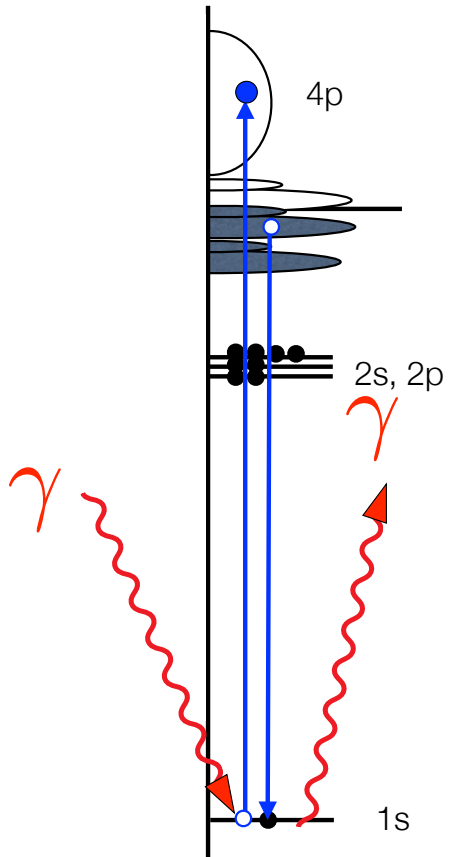
**SIX instrument at NSLS-II (commissioning now)**

*Institute for Theoretical Solid State Physics, IFW Dresden, 01069 Dresden, Germany*

(Received 13 April 2010; published 24 June 2011)



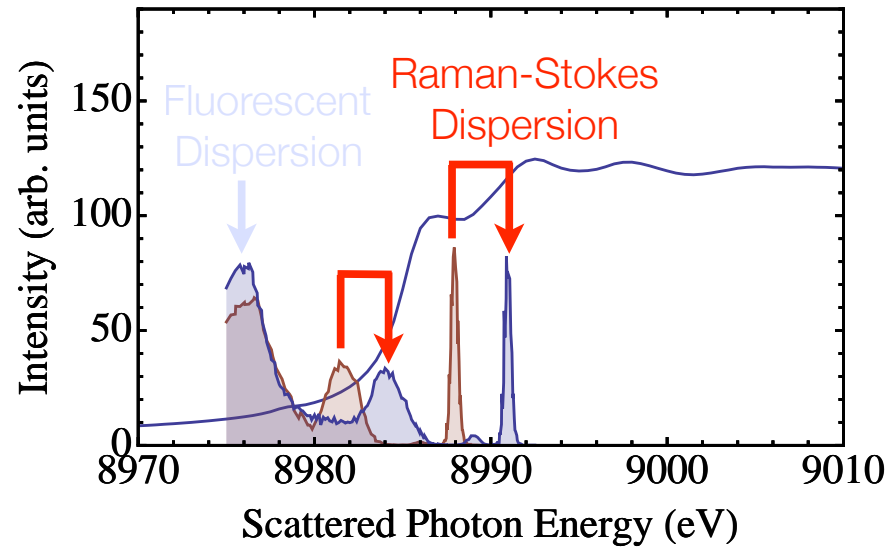
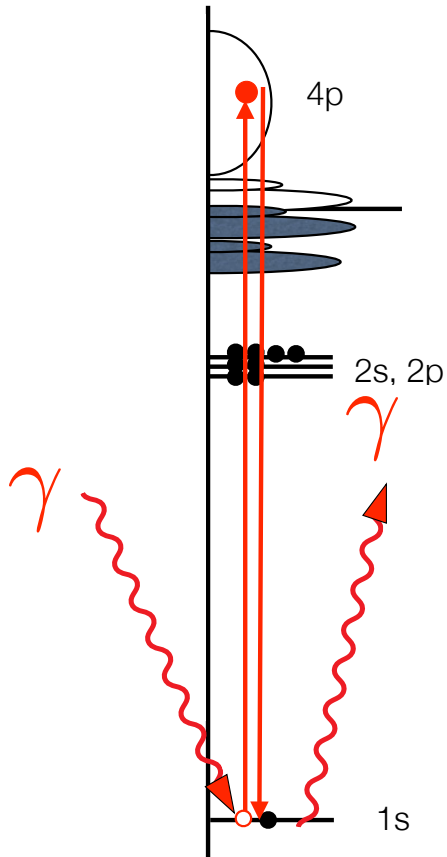
# Energy dispersion in RIXS



Emission line

- Fluorescent-type energy dispersions are independent of  $E_i$

# Energy dispersion

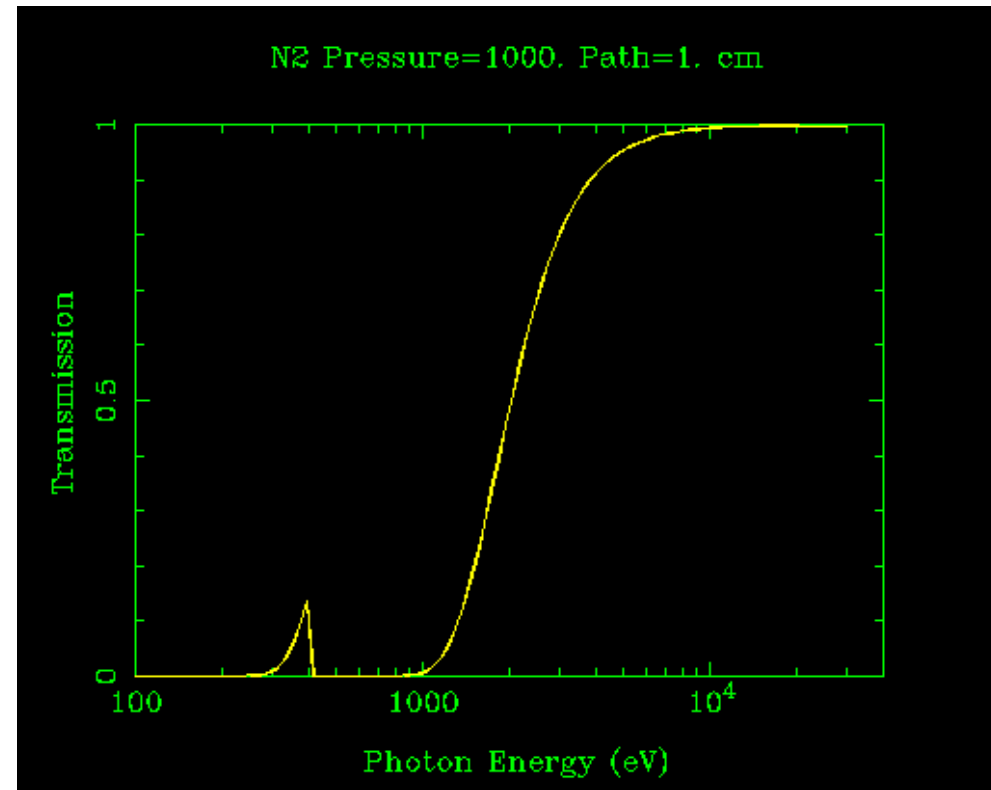
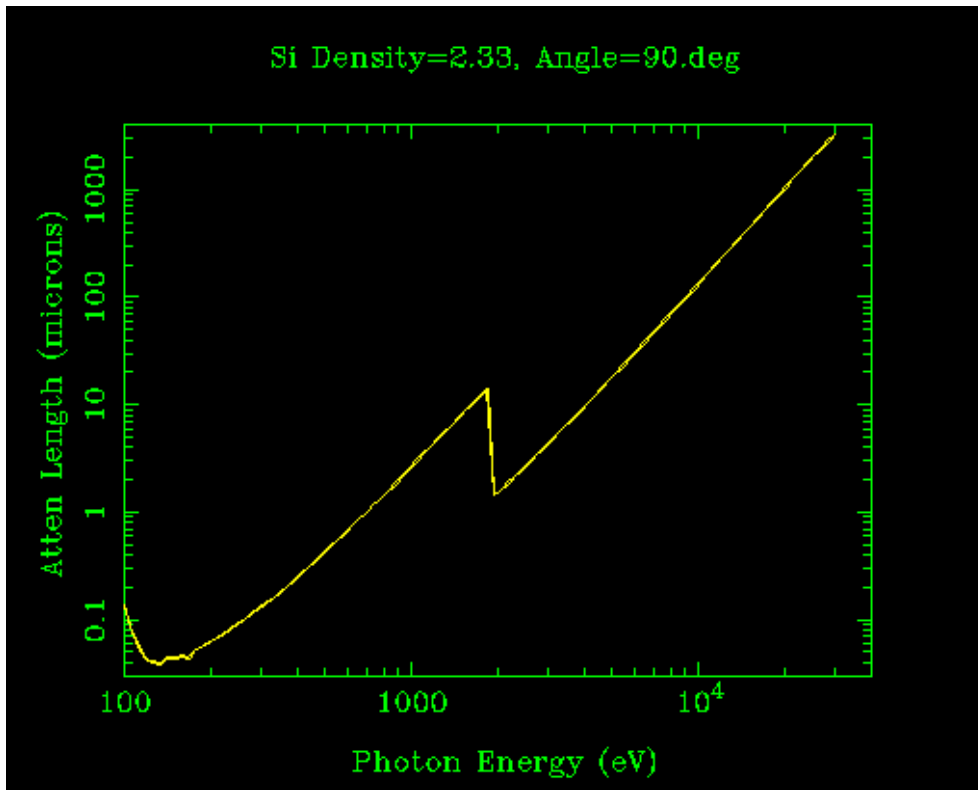


- Features following  $E_i$  represent Raman (RIXS) processes, connected to lower energy (valence) physics

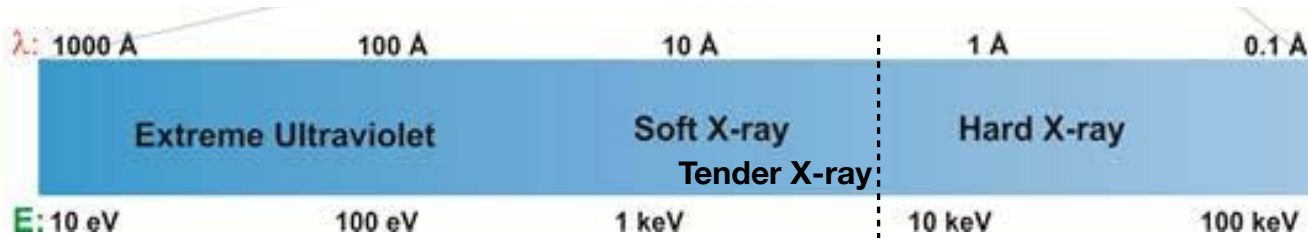
# Practical difference: Soft vs Hard RIXS



[http://henke.lbl.gov/optical\\_constants/](http://henke.lbl.gov/optical_constants/)

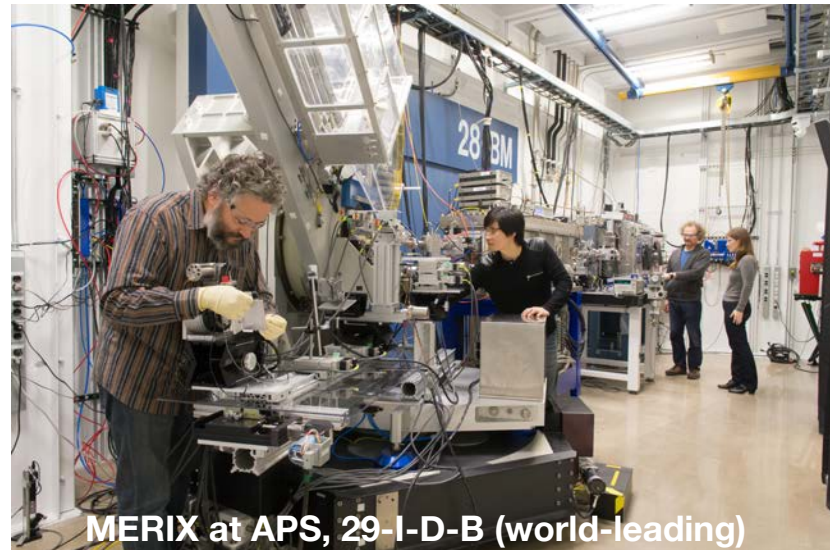


# Practical difference: Soft vs Hard RIXS



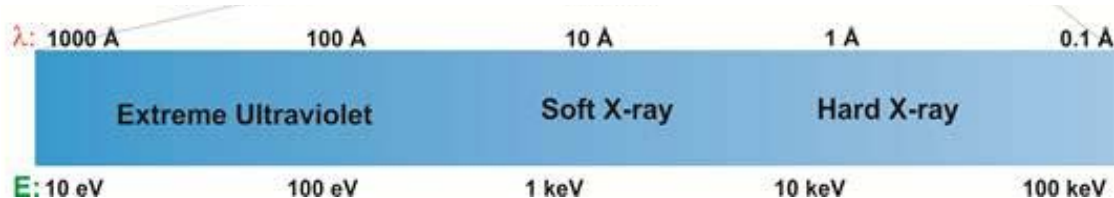
X-rays must be in vacuum  
← More scattering, lower penetration

X-rays can be in helium, air atmosphere  
→ Less scattering, more bulk information

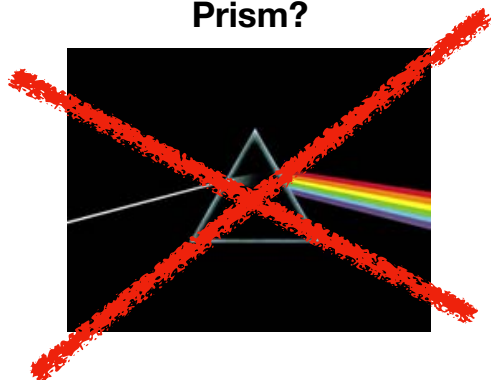


# Technical difference: Soft vs Hard RIXS

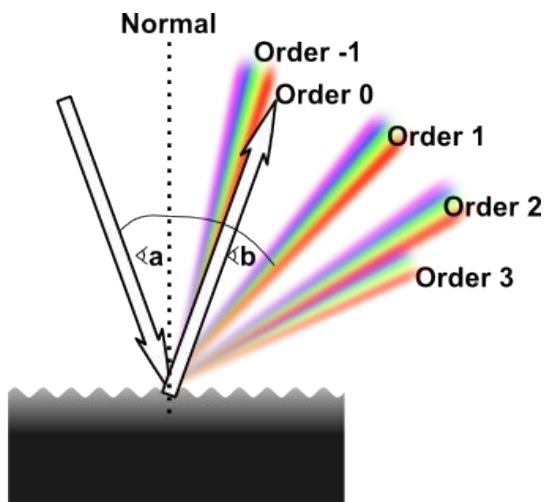
## How to disperse the X-rays?



Prism?



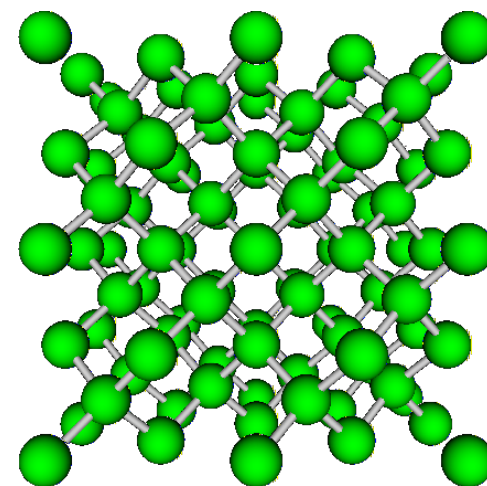
Grating?



If feature sizes can be on order wavelength

Soft X-ray ~ 1 nm

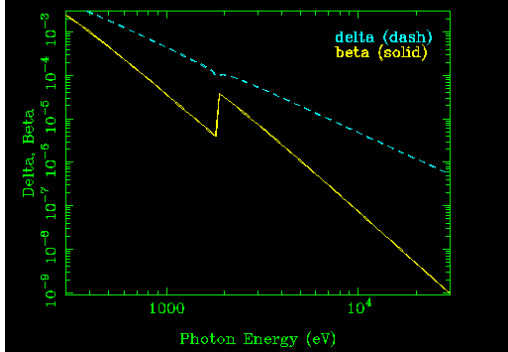
Crystal?



Feature size (atoms) of order X-ray wavelength

Hard X-ray  $\ll$  1nm

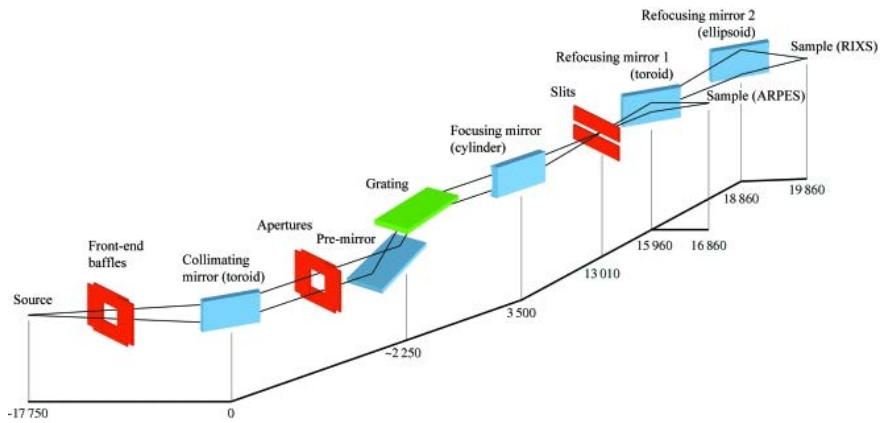
Index of Refraction =  $(1-\delta)-i(\beta)$



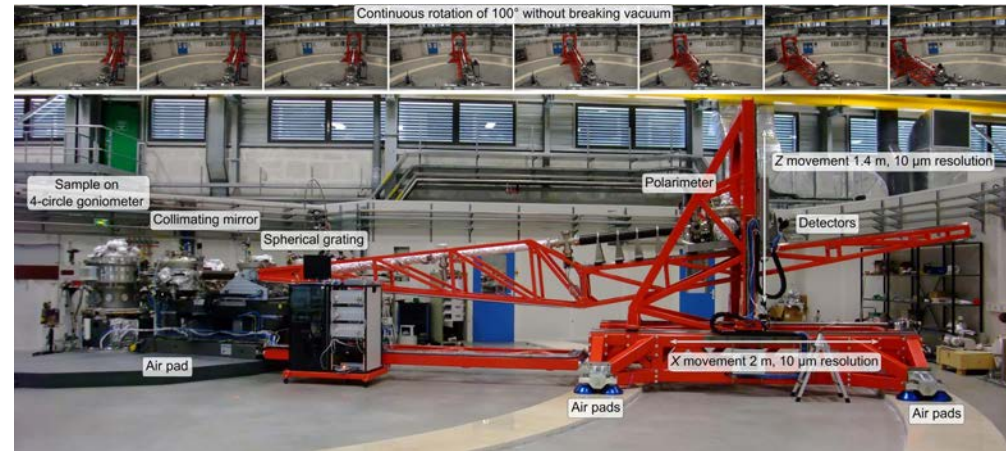
Inefficient

Index changes too small to be effective

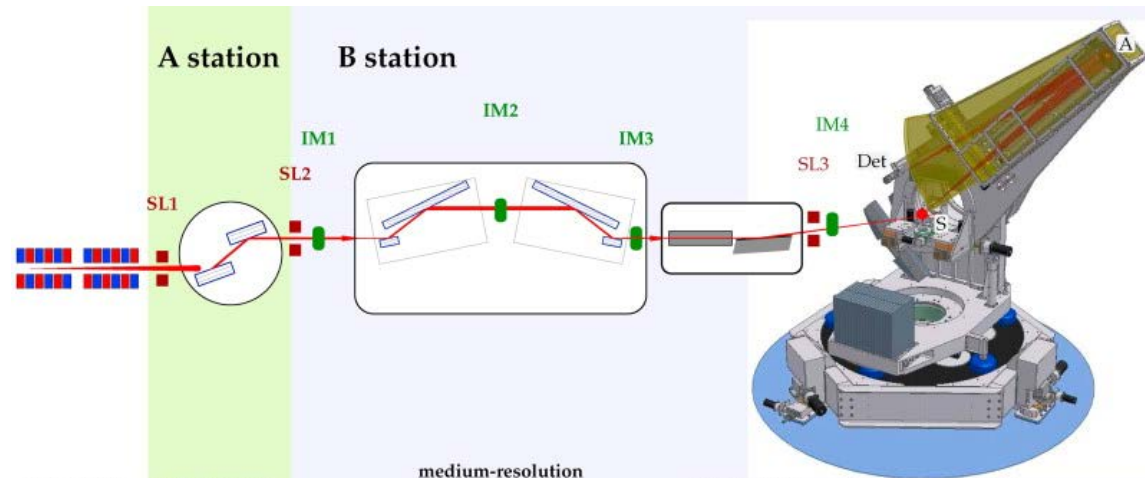
# Technical difference: Soft vs Hard RIXS



**ADDRESS beamline at PSI (no spectrometer)**



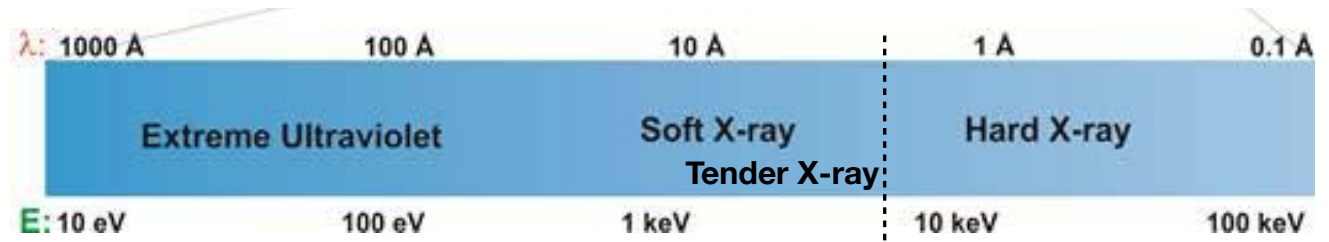
**ERIXS endstation at ESRF (no beamline shown)**



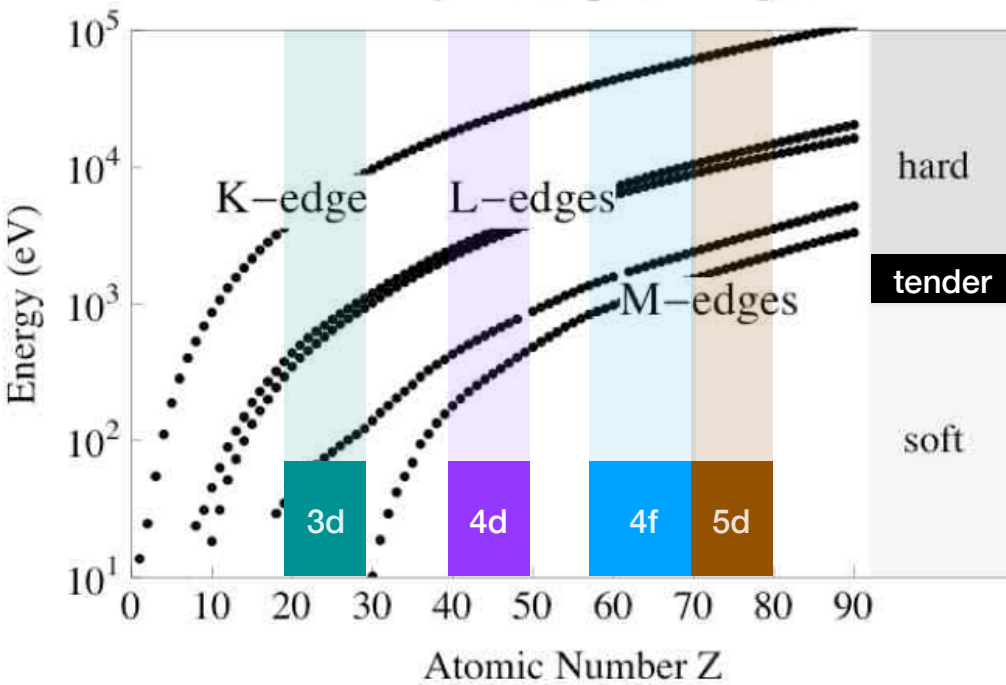
**MERIX spectrometer, sector 27-ID-B, APS**



# Scientific difference: Soft vs Hard RIXS



X-ray Absorption Edges

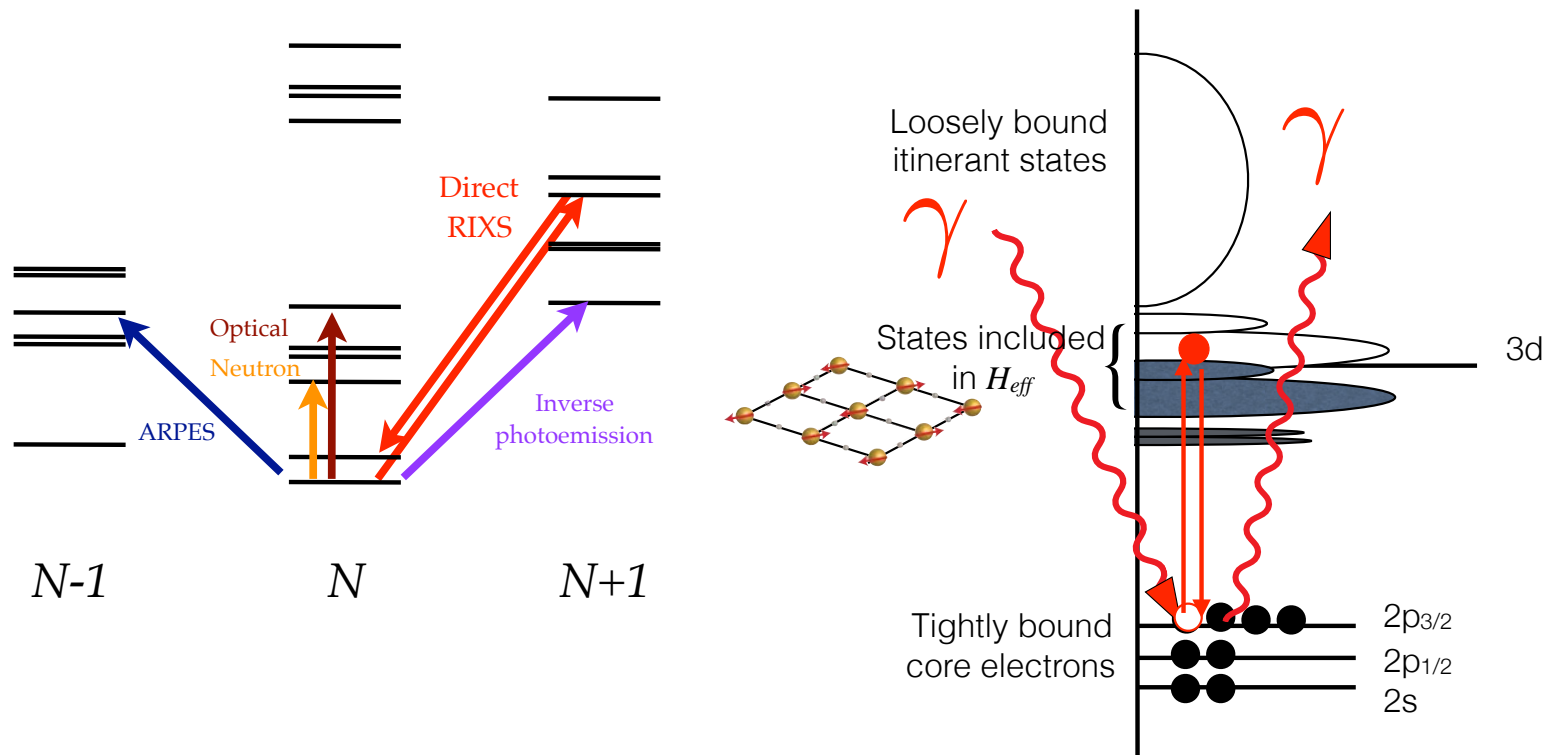


3d transition metal L edges  
Rare-earth/4f M edges

3d transition metal K edges  
Rare-earth/4f L edges  
5d transition metal L edges

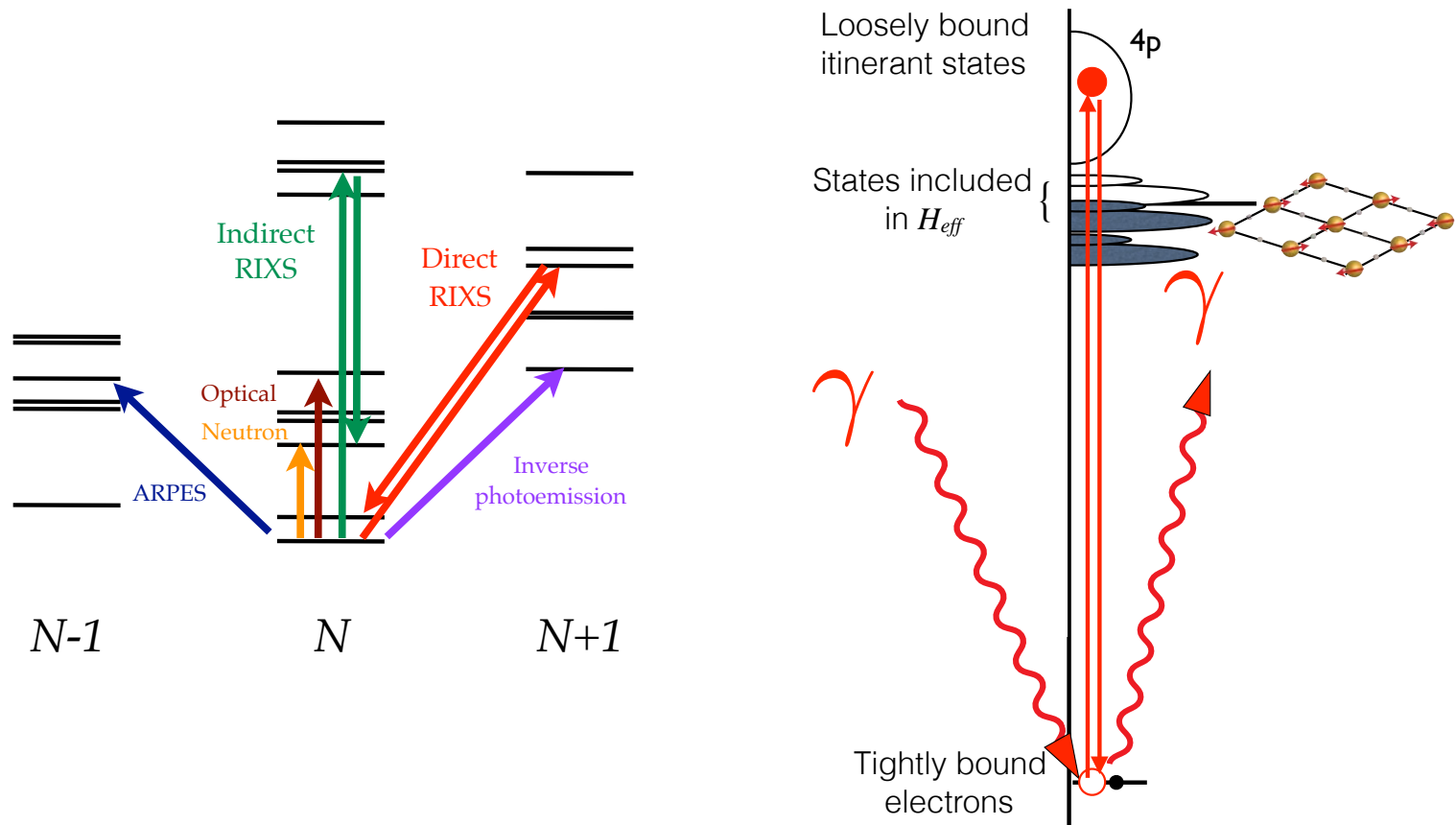
**Edge and energy are determined by elements present and scientific question**

# $N$ -particle spectroscopy



“Direct” RIXS at transition metal  $L$  edges

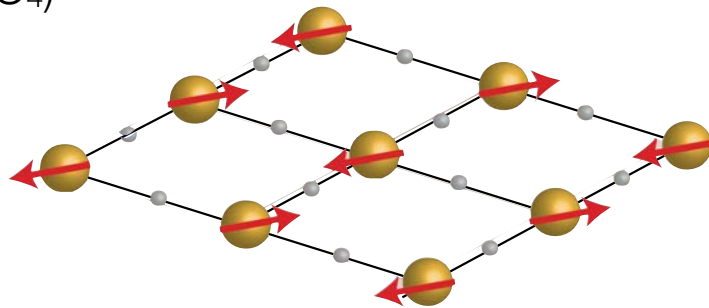
# $N$ -particle spectroscopy



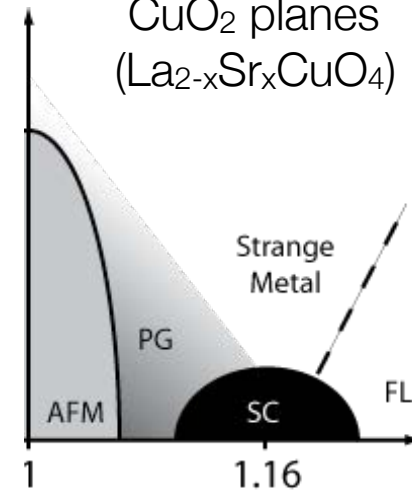
"Indirect" RIXS at transition metal  $K$  edges

# High- $T_c$ cuprate superconductors

1/2-filled  
CuO<sub>2</sub> planes  
(La<sub>2</sub>CuO<sub>4</sub>)



Hole-doped  
CuO<sub>2</sub> planes  
(La<sub>2-x</sub>Sr<sub>x</sub>CuO<sub>4</sub>)

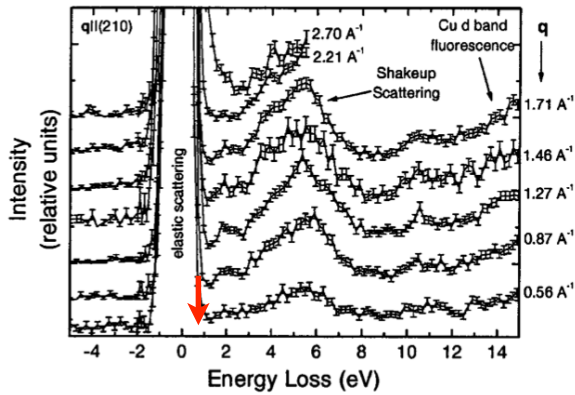


- Strong repulsion favors antiferromagnetic state
- Hubbard model describes AFM

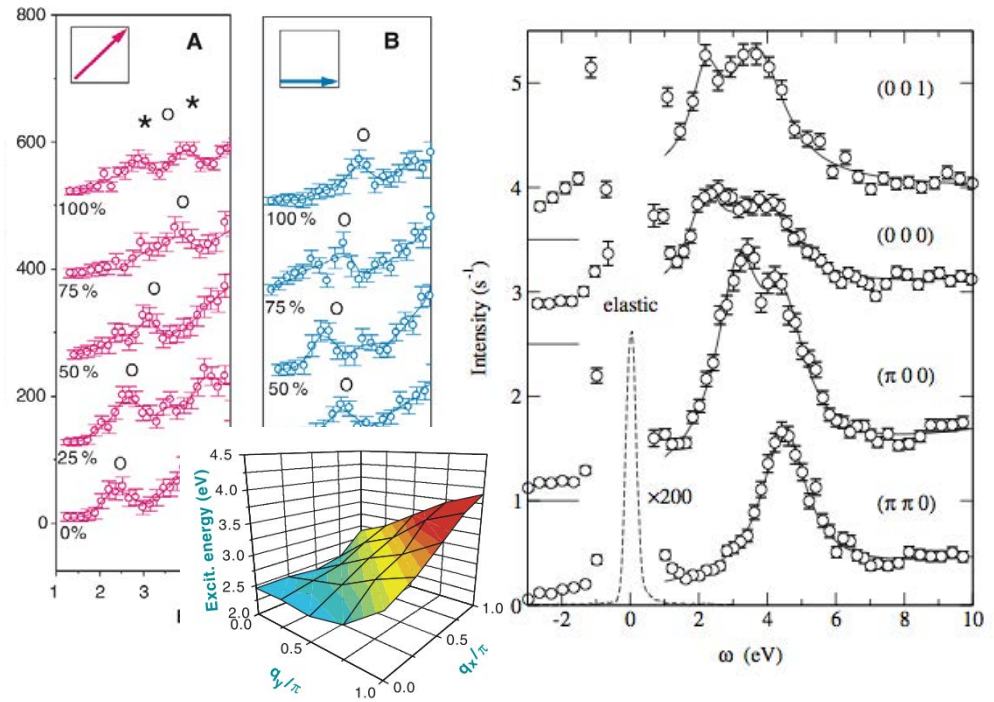
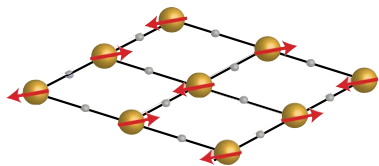
Hole doping a 2D Mott insulator induces high-temperature superconductivity

$$H = -t \sum_{\langle i,j \rangle, \sigma} (c_{i,\sigma}^\dagger c_{j,\sigma} + h.c.) + U \sum_{i=1}^N n_{i\uparrow} n_{i\downarrow}$$

# Mott gap dispersion in high-Tc superconductors (Cu K edge)



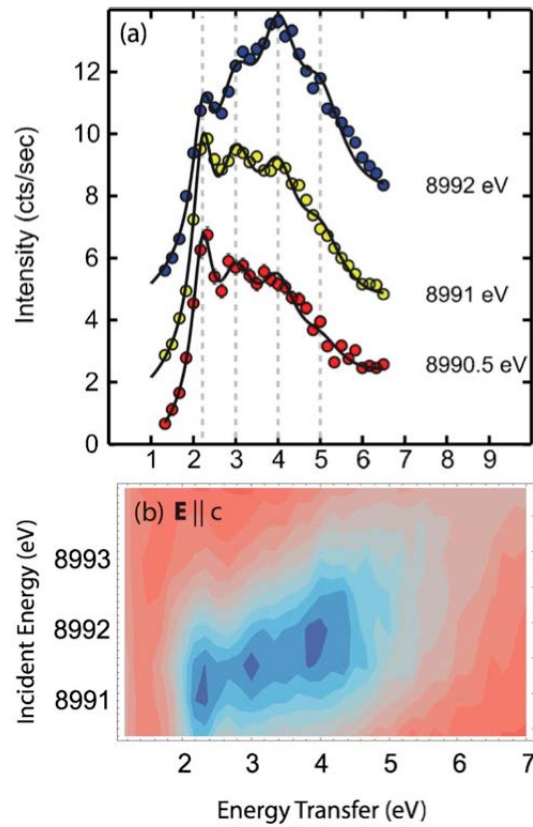
Abbamonte *et al.* PRL 83, 860 (1999)



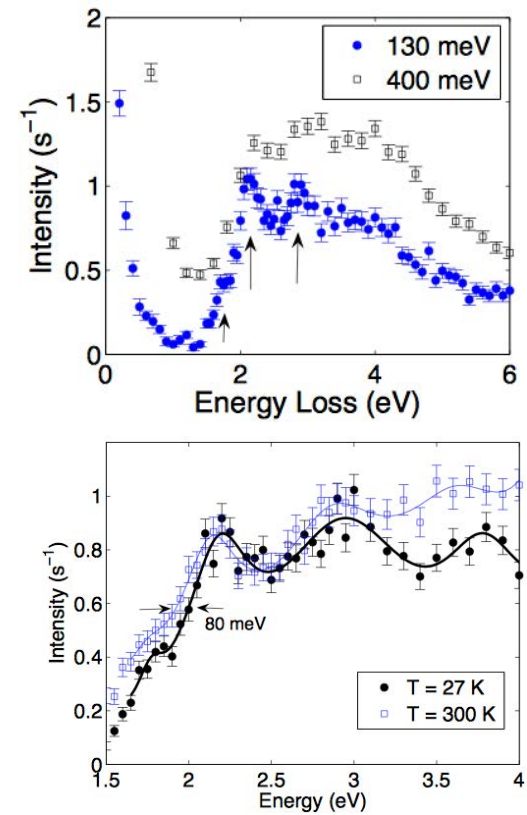
Hasan *et al.* Science 288, 1811 (2000)

Kim *et al.* PRL 89, 177003 (2002)

# Incident energy and temperature dependence

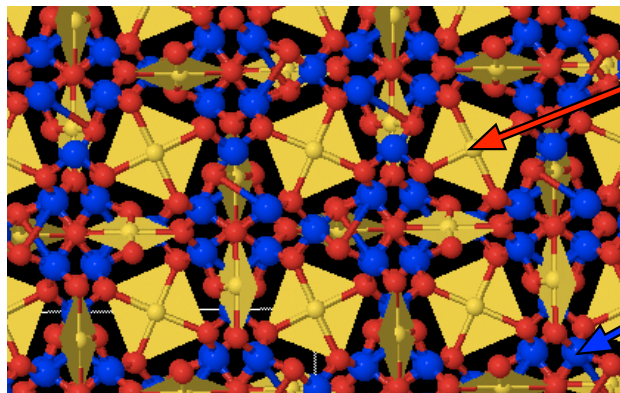


Lu, JNH, *et al.* *PRL* **95**, 217003 (2005)  
Lu, JNH, *et al.* *PRB* **74**, 224509 (2006)



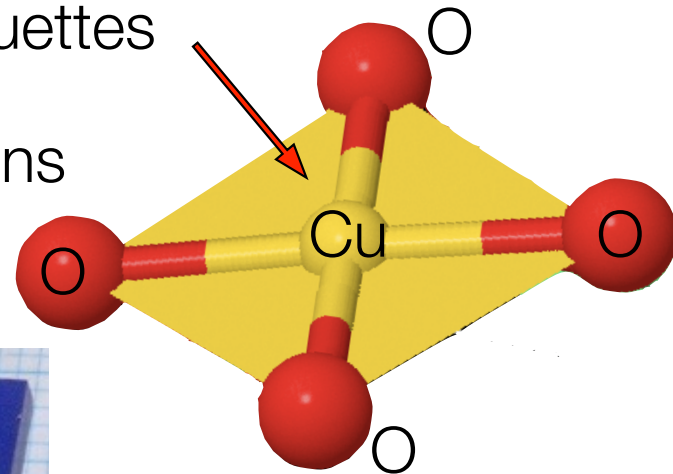
Ellis, *et al.* *PRB* **77**, 060501 (2008)

# Experimental realization of 0-D HTSC



CuB<sub>2</sub>O<sub>4</sub>, <001> view

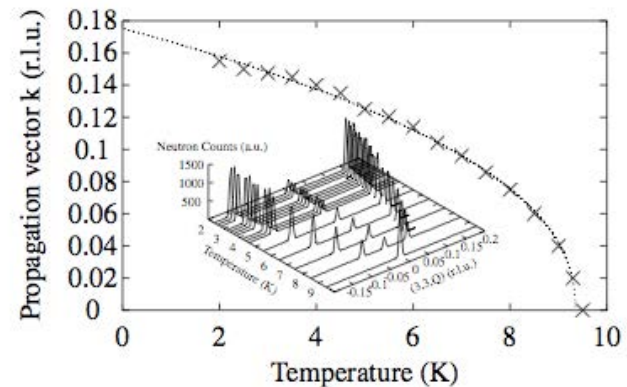
CuO<sub>4</sub>  
plaquettes  
B<sup>3+</sup> ions



- Low-Z nonresonant species
- Tractable electronic structure
- Ideal system for exploring RIXS phenomenology

Hancock *et al.* *PRB* **80**, 092509 (2009)

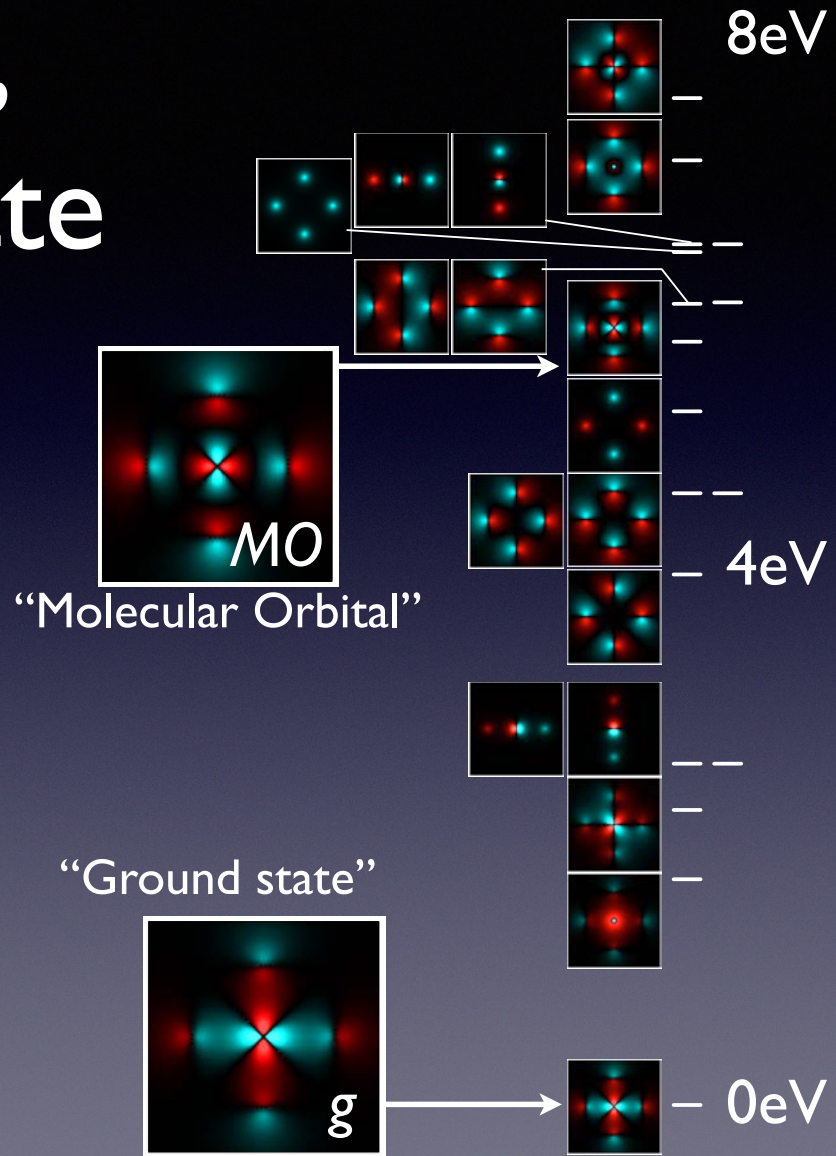
Hancock *et al.* *NJP* **12**, 033001 (2010)



B. Roessli *et al.*, *PRL* **86**, 1885 (2001)

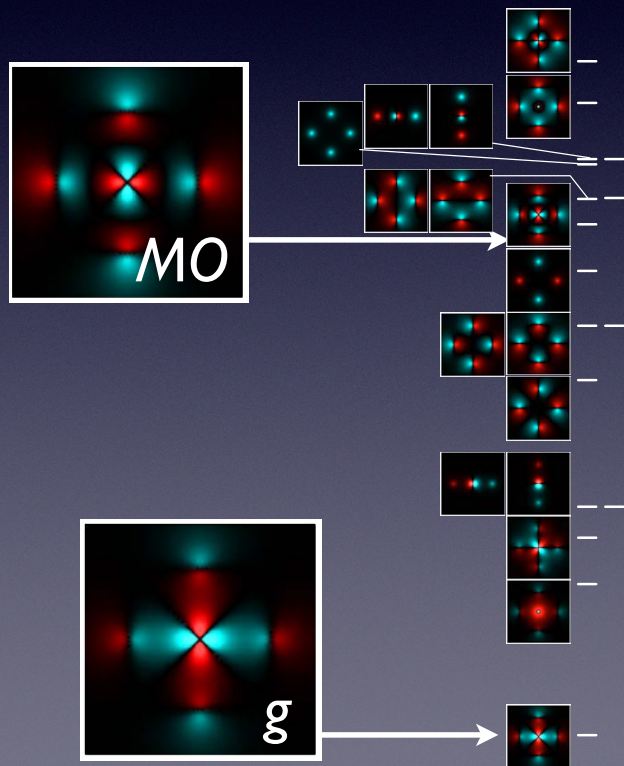
# One hole, one plaquette

- Many wavefunctions of varied angular symmetry
- Ground and molecular orbital states are the only ones of  $x^2-y^2$  symmetry
- Simple RIXS selection rule (for indirect,  $K$  edge transitions): symmetry=constant

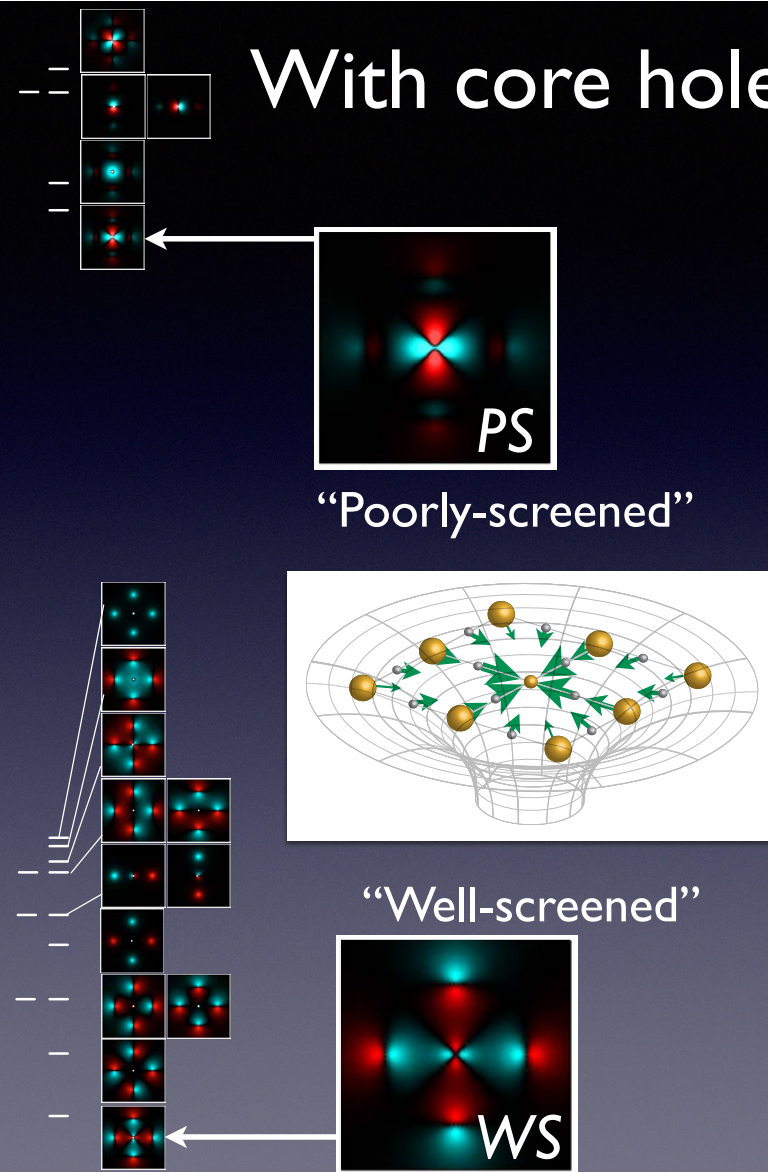




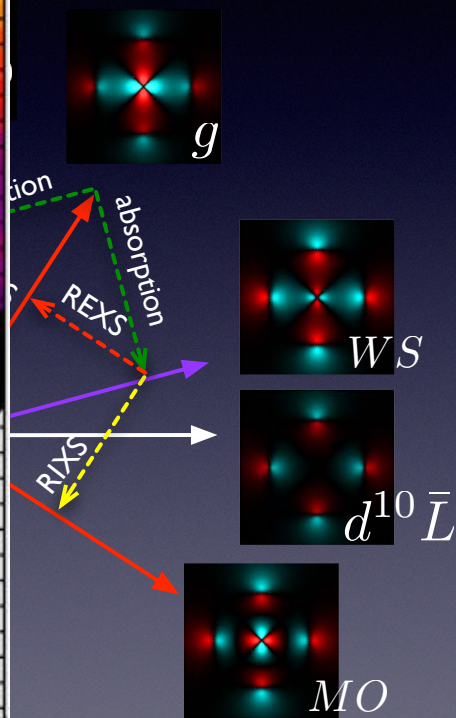
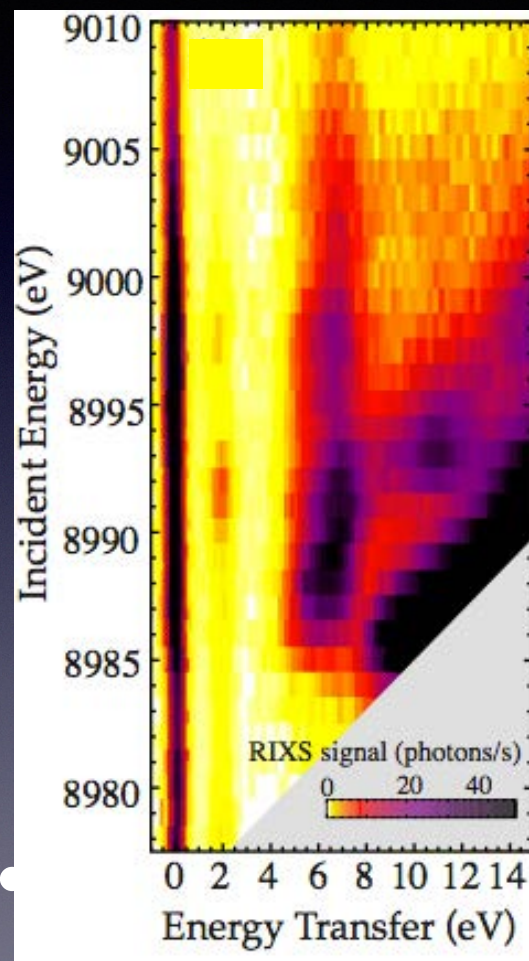
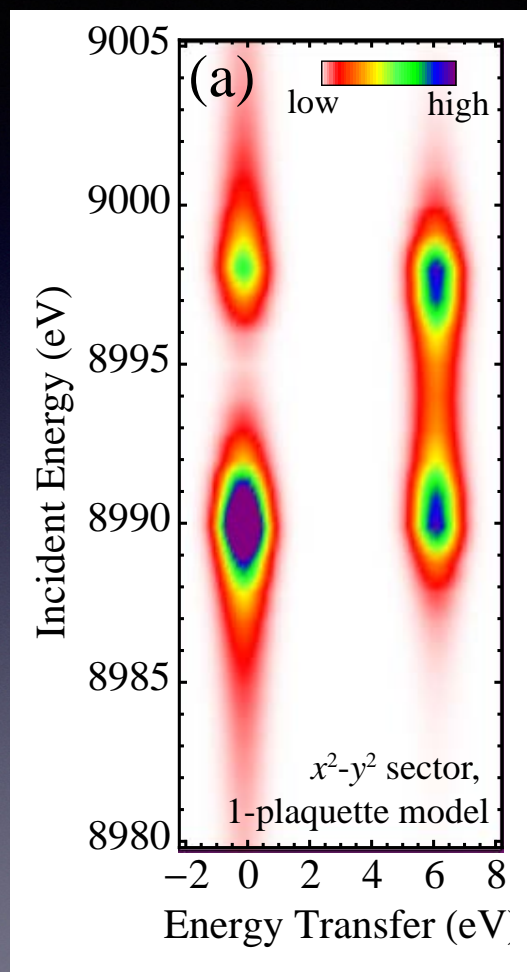
# Without core hole



# With core hole



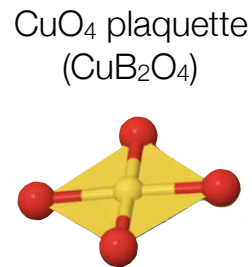
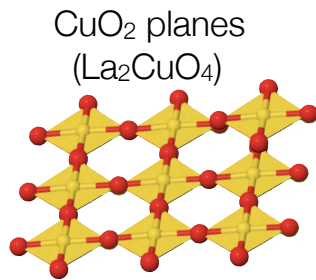
# I-plaquette RIXS calculation



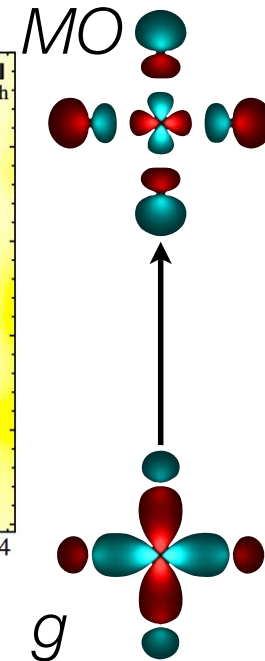
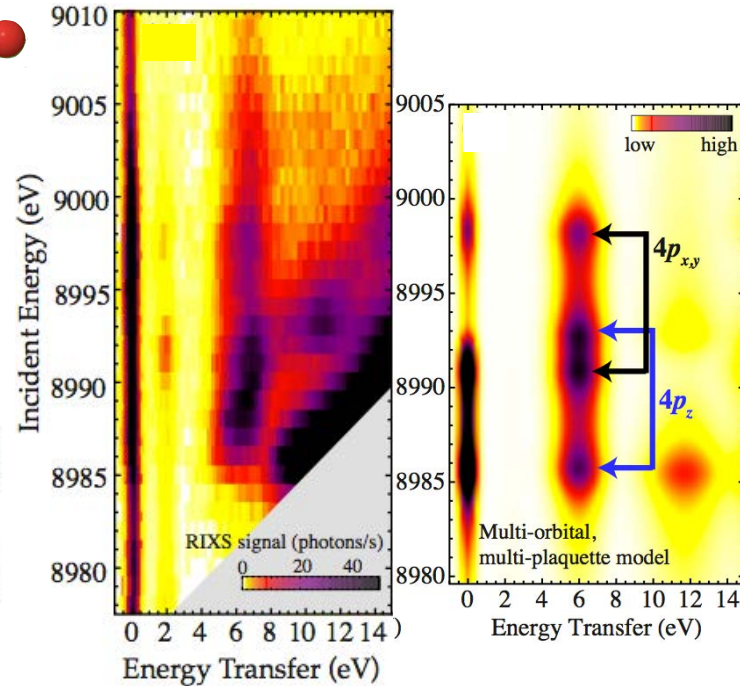
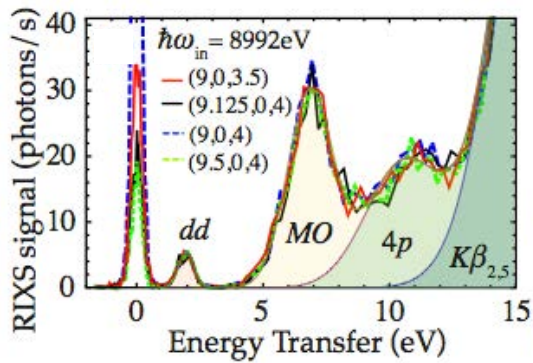
capable of giving

phase information

# A simple cuprate, $\text{CuB}_2\text{O}_4$



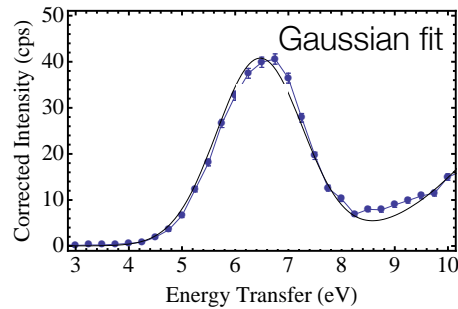
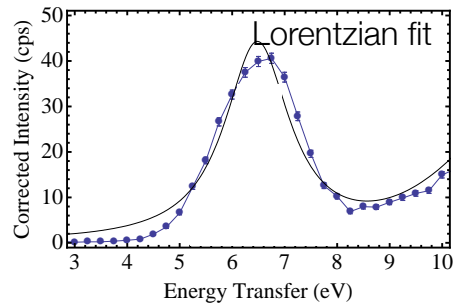
- Cluster modeling gave good, but imperfect agreement



Hancock *et al.* *PRB* **80**, 092509 (2009)

Hancock *et al.* *NJP* **12**, 033001 (2010)

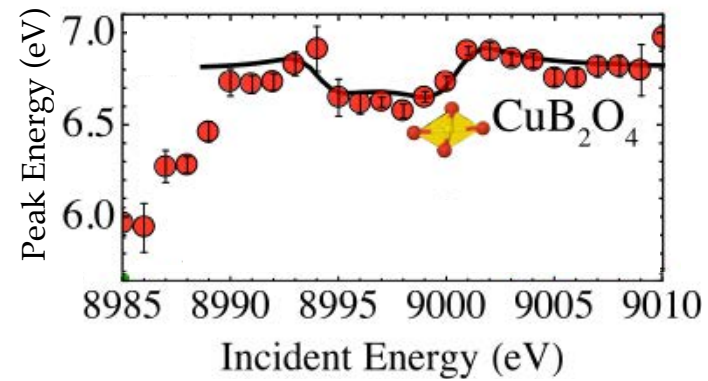
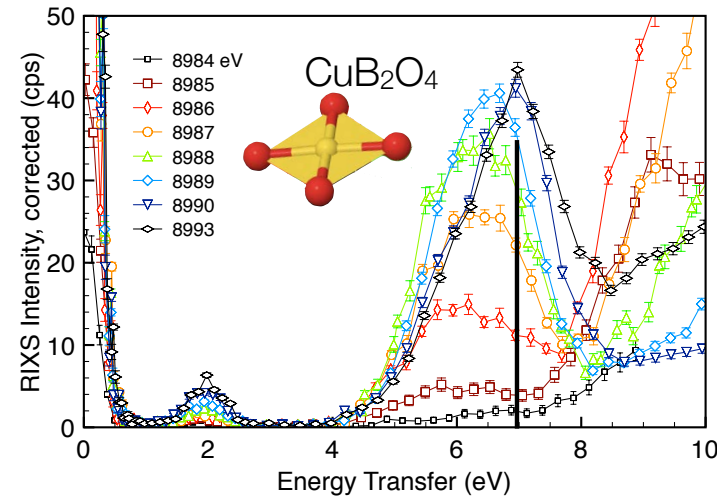
# Elusive properties of MO excitation



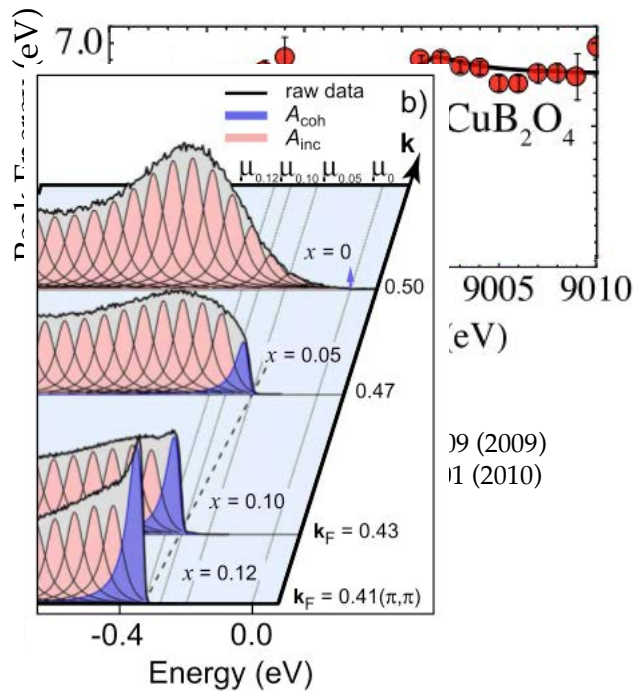
- Gaussian lineshape
- Shifting resonance behavior

Hancock *et al.* *PRB* **80**, 092509 (2009)

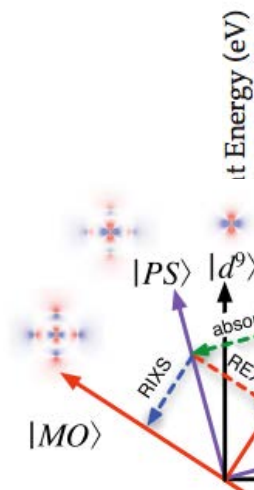
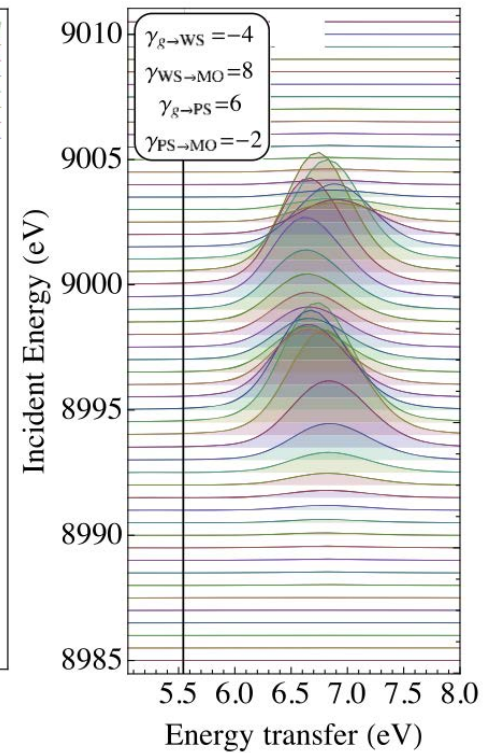
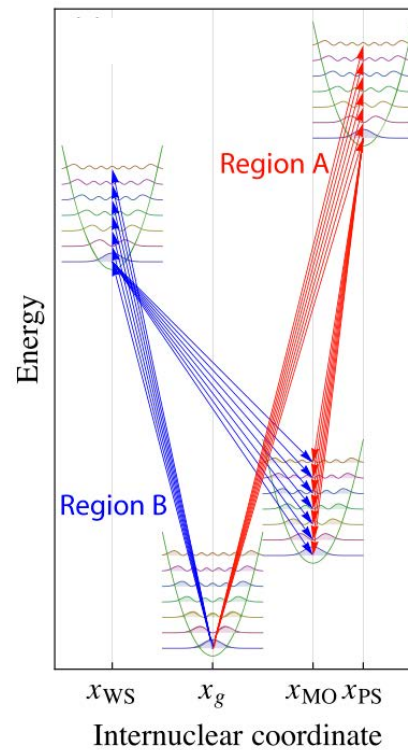
Hancock *et al.* *NJP* **12**, 033001 (2010)



# Electron-phonon coupling at the MO excitation



Shen, *et al*, PRL **93**, 267002 (2004)

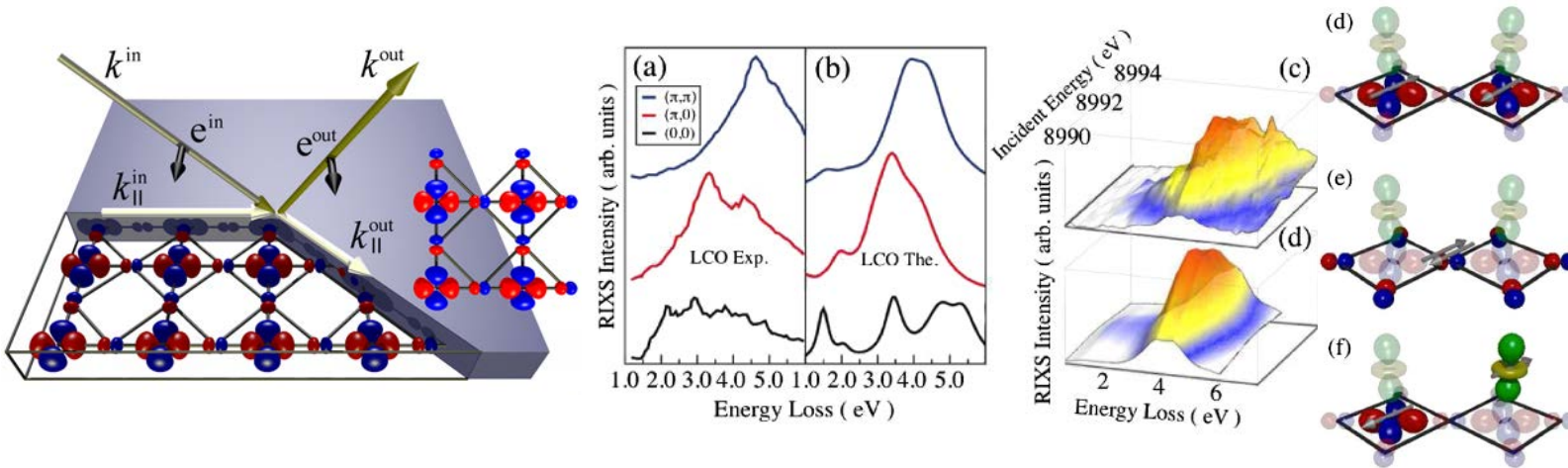


- Simple Franck-Condon model explains remaining features of the spectra
- Gives an estimate of the electron-phonon coupling parameter

# Theory of RIXS in the cuprates

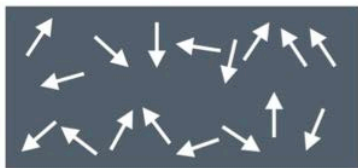
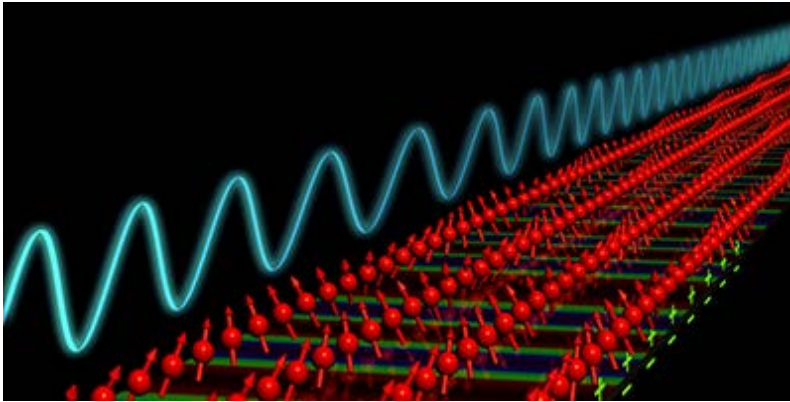
## Kramers-Heisenberg Formula

$$I(\omega, \mathbf{q}, \omega', \mathbf{q}') = \sum_f \left| \sum_n \frac{\langle f|T|n\rangle \langle n|T|g\rangle}{\hbar\omega - E_n + i\frac{\Gamma}{2}} \right|^2 \delta(\hbar\omega' + E_f - \hbar\omega - E_g),$$



- Calculations model RIXS cross section to understand electronic structure of high- $T_c$  superconductors

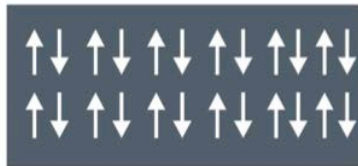
# Magnetically-ordered states



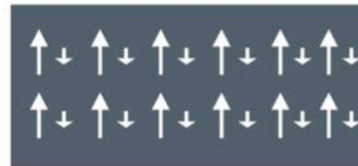
A = paramagnetic



B = ferromagnetic



C = antiferromagnetic

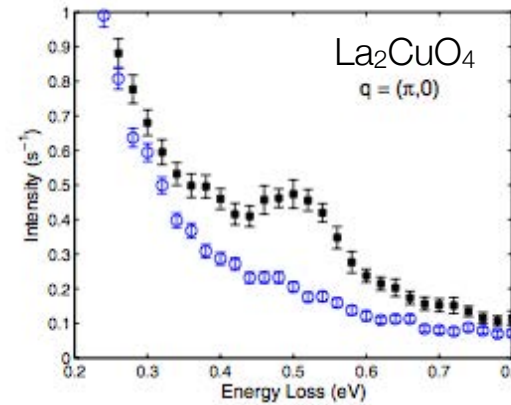


D = ferrimagnetic

- Magnetic interactions between magnetic ions can lead to ordered states
- Disturbance in order forms “spin waves” or “magnons”
- Like sound, magnons carry momentum and energy
- Dispersion relation is important

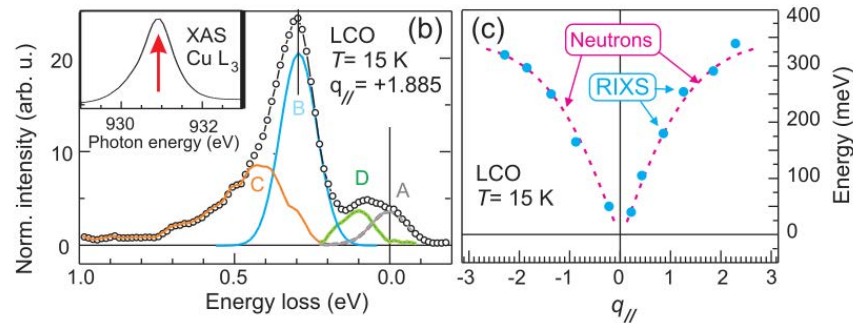
# Magnetic excitations via RIXS

- 2008 Cu K edge RIXS showed bi-magnon feature, consistent with Raman  $q=0$  result



Hill, Blumberg, et al, PRL 100 097001 (2008)

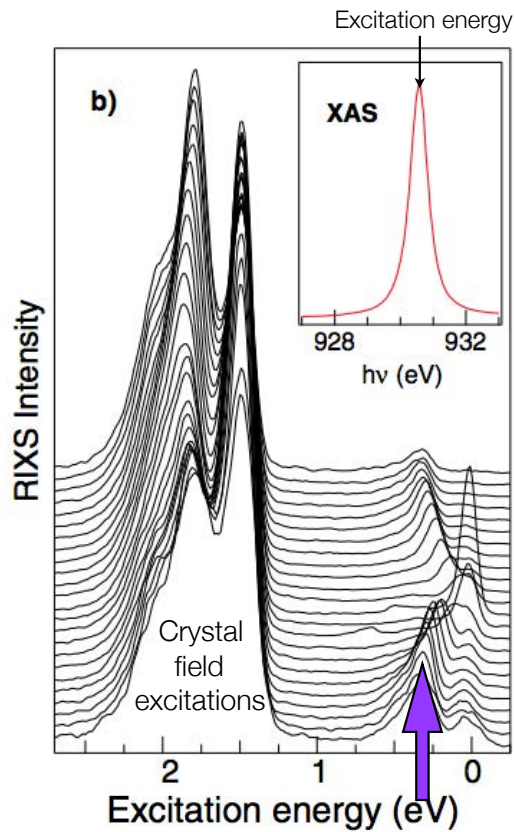
- 2009 Cu L edge (soft X-ray) RIXS showed a shoulder of elastic line consistent with *single*-magnon excitation



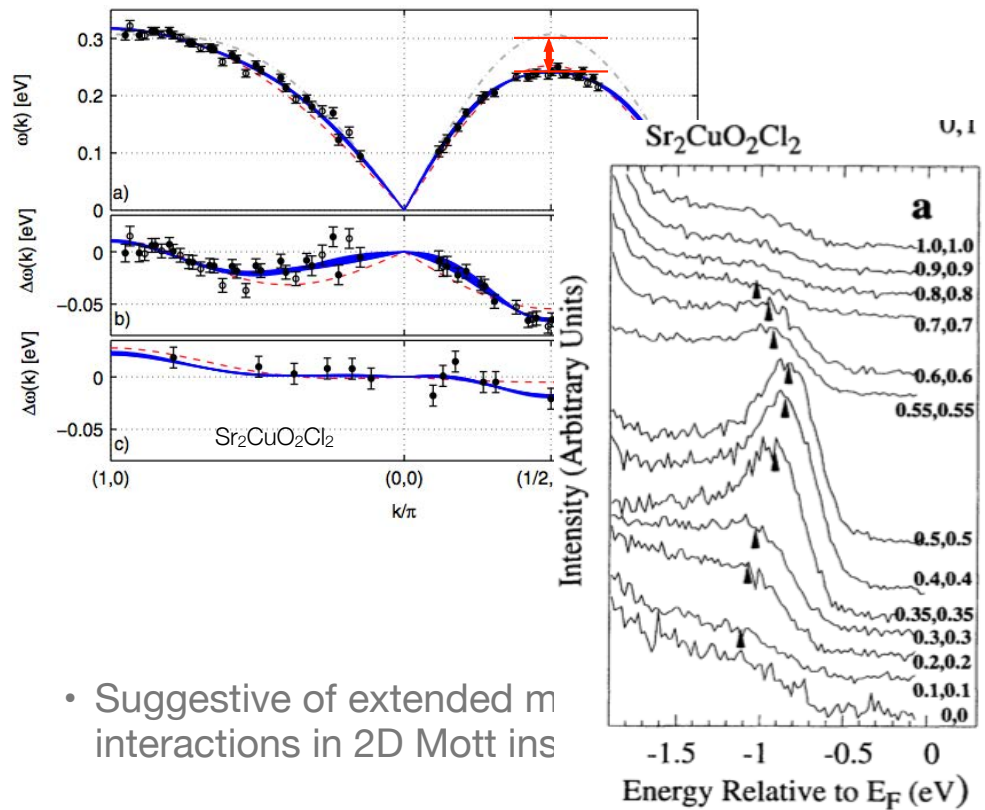
Braicovich, Ghiringhelli, et al, PRL 102 167401 (2009)



# Extended interactions in a Mott insulator $\text{Sr}_2\text{CuO}_2\text{Cl}_2$



Guarise, ...JNH, ... *et al.* PRL 12, 033001 (2010)

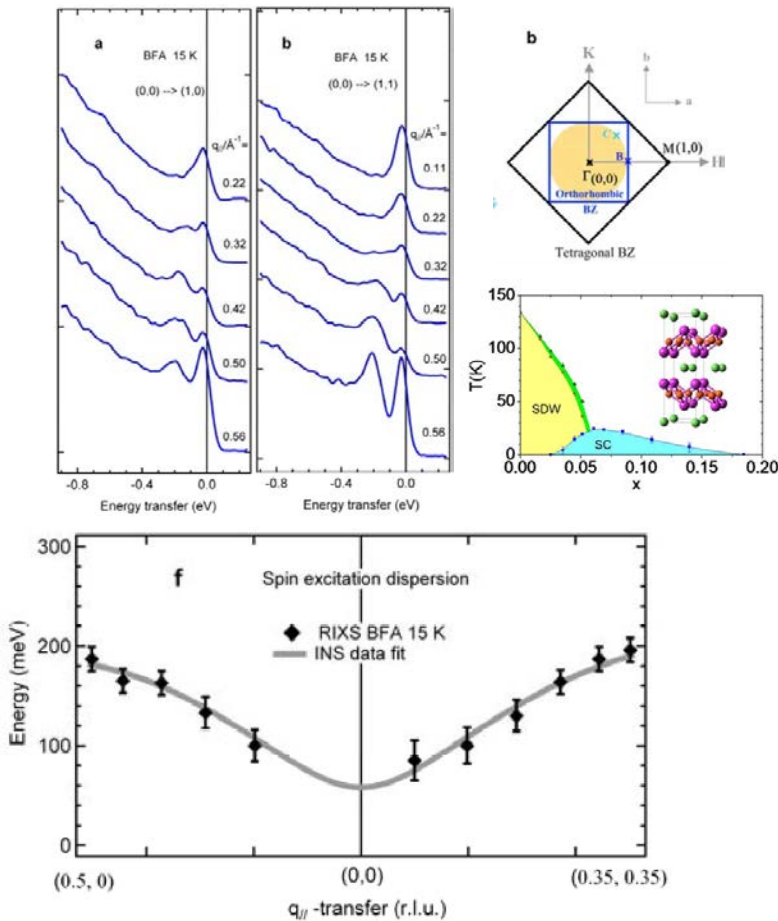


- Suggestive of extended m interactions in 2D Mott ins

Wells, *et al.* PRL 74, 964 (1994)

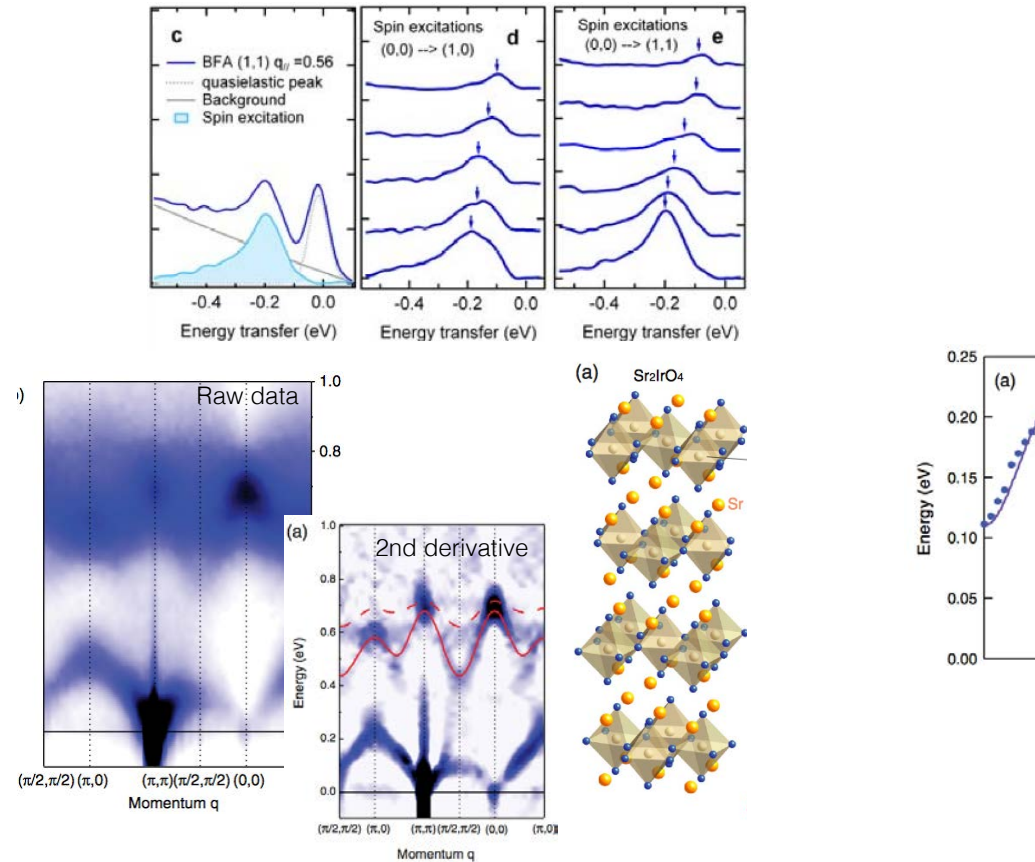
# Iron pnictide superconductors and Correlated 5d systems

## BaFe<sub>2</sub>As<sub>2</sub>



Zhou, et al, *Nature Comm.* **4**, 1470 (2013)

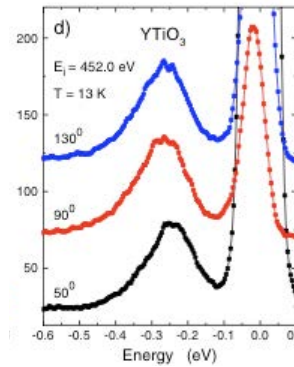
## SrIr<sub>2</sub>O<sub>4</sub>



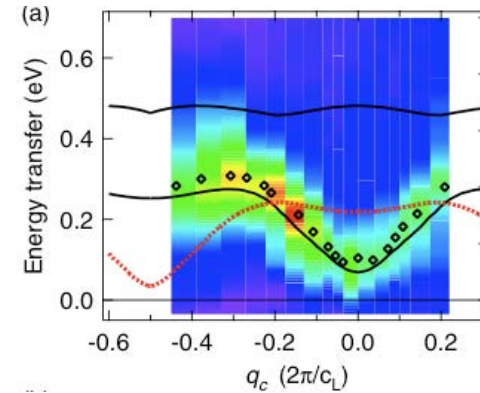
Kim, et al, *Phys. Rev. Lett.* **108**, 177003 (2012)

# Other RIXS highlights

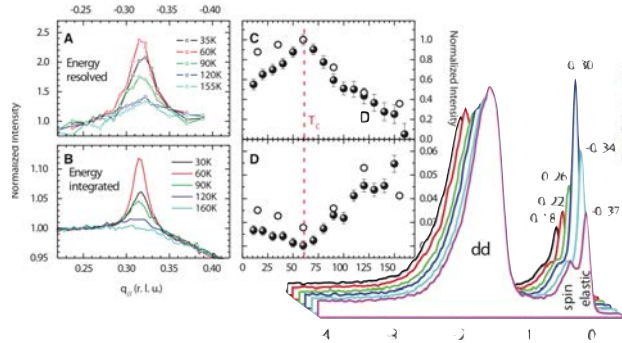
- **Orbiton** dispersion
- **Phonons** excited through core hole intermediates
- **Spin-2** (triplon) excitations
- **Incipient CDW** effect at  $T_c$  in YBCO
- **Paramagnon** fluctuations in doped cuprates



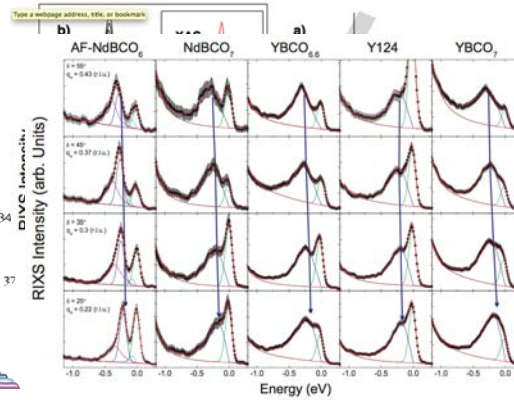
Ulrich, et al PRL 103, 107205 (2009)



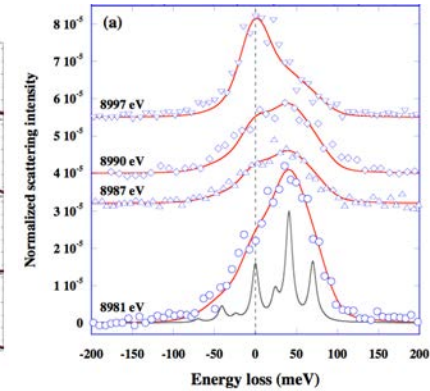
Schlappa, et al PRL 103, 047401 (2009)



Ghiringhelli, et al Science 337, 821 (2012)



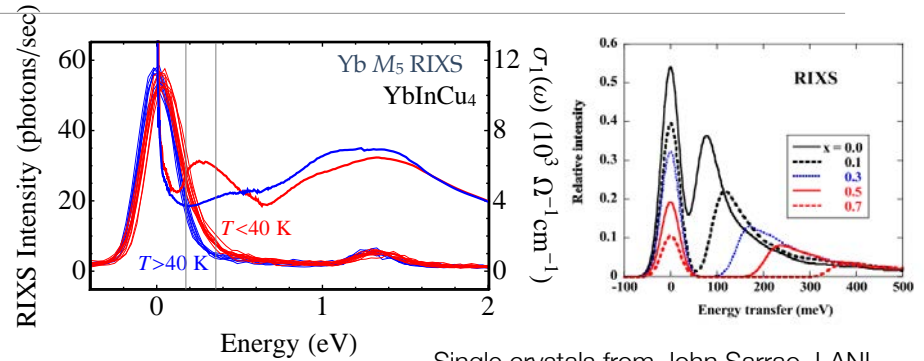
Le Tacon, et al, Nat. Phys. 2011



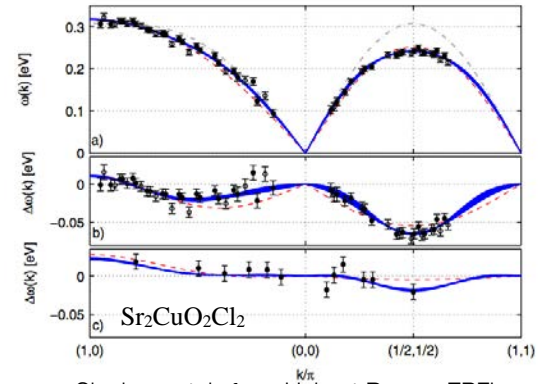
Yavas, et al JPCM 22, 485601 (2010)

# Soft X-ray RIXS work

- Rediscovery of the Kondo gap (now with X-rays)
- 1st mapping of magnon spectra in  $\text{Sr}_2\text{CuO}_2\text{Cl}_2$  (non-local exchange)
- Connected RIXS to edge singularity physics



Single crystals from John Sarrao, LANL



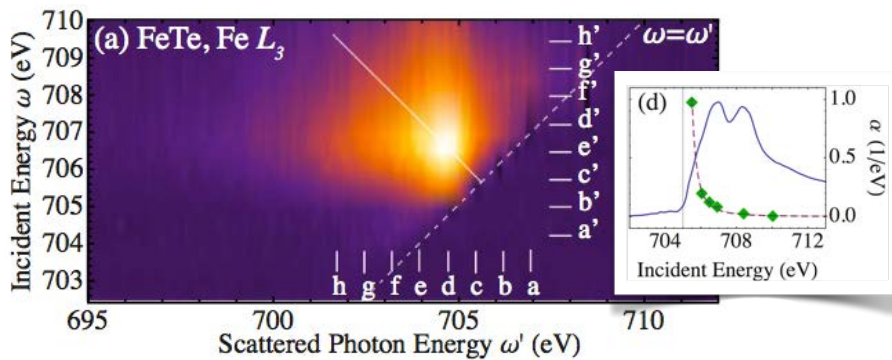
Single crystals from Helmut Berger, EPFL

Hancock *et al.* (2011)

Kotani, in review (2011)

Guarise *et al.* PRL **12**, 033001 (2010)

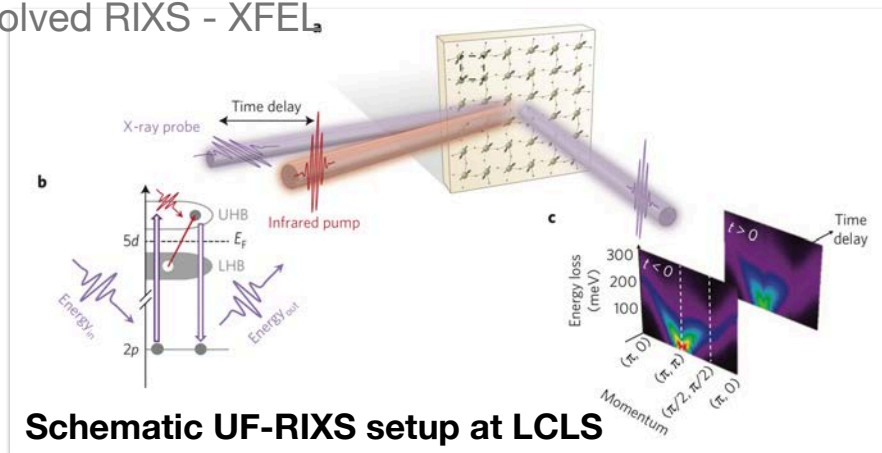
Hancock *et al.* PRB **82**, 020513R (2010)



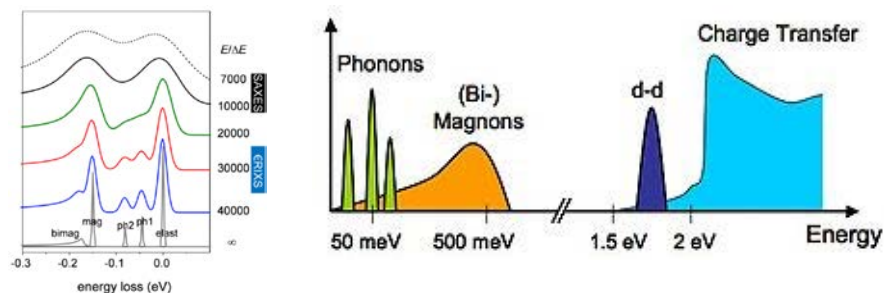
Single crystals from Enrico Giannini group, Geneva

# Future technique directions with RIXS

- Time-resolved RIXS - XFELs



- Higher energy resolution

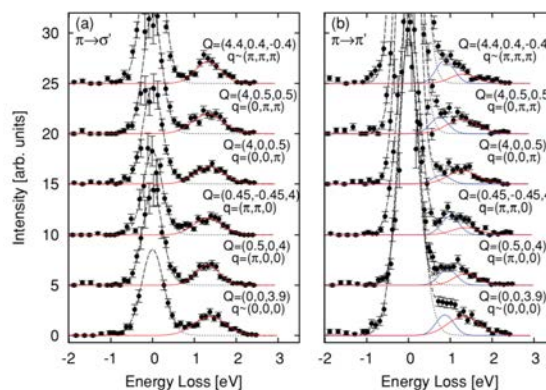


**Energy scales in correlated electron systems**

- Tender X-ray RIXS - new edges and opportunities



- Polarization analysis of scattered photon



**High potential to resolve subtle changes, identify excitation channels**

# Summary/outlook

---

- RIXS is a 21st century technique, deep probe of exotic excitations
- Many opportunities on horizon for new generation of synchrotron scientists
- APS and NSLS-II have world-leading opportunities in hard and soft RIXS science
  
- New opportunities for science of correlated electron materials
  - Mixed-valent/Kondo lattice physics
  - High-temperature superconductivity
  - Search for exotic phases for quantum information science

**Thank you!**